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PRODUCTION OF W^+W^- AND $t\bar{t}$ PAIRS VIA PHOTON-PHOTON MECHANISM IN PROTON-PROTON SCATTERING

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Abstract

We review our recent results for production of W^+W^- and $t\bar{t}$ pairs via photon-photon fusion mechanism. A sketch of theoretical approach is presented. We include the transverse momenta of the photons in the calculation of photon fluxess. Then we present our results for the cross section (total and differential) of W^+W^- production. Results for different parametrizations of proton structure functions are used to calculate the inelastic fluxes of photons. A discussion on the rapidity gap survival probability, due to remnant fragmentation, is presented. A similar discussion is presented for $t\bar{t}$ production.

1 Introduction

It was realized rather recently that the electroweak corrections are important for precise calculations of cross sections in different processes. The $pp \to W^+W^-$ process is a good example (see e.g. [1]) and $\gamma\gamma \to W^+W^-$ is a relevant subprocess. This subprocess is important also in the context of searches beyond the Standard Model [2, 3]. By imposing special conditions on the final state, this contribution can be observed experimentally [4, 5].

In [6, 7] we developed a formalism tp calculate $pp \to l^+l^-$ processes, proceeding via photon-photon fusion. In [8] we used the same technique to calculate the cross section for $pp \to W^+W^-$ reaction proceeding via photon-photon fusion. In order to make reference to real "measurements" of the photon-photon contribution one has to include in addition the gap survival probability caused by extra emissions. In [9] we concentrated on the effect related to remnant fragmentation and its destroying of the rapidity gap.

In [10] we calculated the cross section for the photon-photon contribution to the $pp \to t\bar{t}$ reaction, including also the effects of gap survival probability.

Here we briefly review our results obtained in [8, 9, 10].

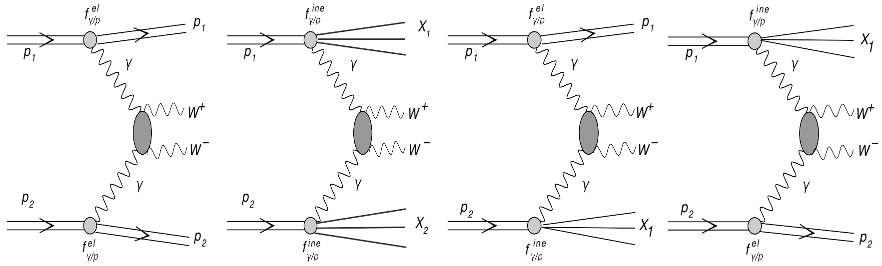

Figure 1: *Diagrams representing different types of photon-photon induced mechanisms for production of W^+W^- pairs.*

2 A skech of the formalism

In our analyses we included the different types of processes shown in Fig. 1.

In our approach we include transverse momenta of (virtual) photons. Then the differential cross section for W^+W^- production can be written as:

$$\frac{d\sigma^{(i,j)}}{dy_1 dy_2 d^2\vec{p}_{T1} d^2\vec{p}_{T2}} = \int \frac{d^2\vec{q}_{T1}}{\pi \vec{q}_{T1}^{\,2}} \frac{d^2\vec{q}_{T2}}{\pi \vec{q}_{T2}^{\,2}} \mathcal{F}^{(i)}_{\gamma^*/A}(x_1, \vec{q}_{T1})\, \mathcal{F}^{(j)}_{\gamma^*/B}(x_2, \vec{q}_{T2}) \frac{d\sigma^*(p_1, p_2; \vec{q}_{T1}, \vec{q}_{T2})}{dy_1 dy_2 d^2\vec{p}_{T1} d^2\vec{p}_{T2}}\,, \quad (1)$$

where $i, j =$ elastic, inelastic and the longitudinal momentum fractions are expressed in terms of rapidities and transverse momenta of W bosons.

The elementary off-shell cross section in (1) is written as:

$$\frac{d\sigma^*(p_1, p_2; \vec{q}_{T1}, \vec{q}_{T2})}{dy_1 dy_2 d^2\vec{p}_{T1} d^2\vec{p}_{T2}} = \frac{1}{16\pi^2 (x_1 x_2 s)^2} \sum_{\lambda_{W^+} \lambda_{W^-}} |M(\lambda_{W^+}, \lambda_{W^-})|^2 \, \delta^{(2)}(\vec{p}_{T1} + \vec{p}_{T2} - \vec{q}_{T1} - \vec{q}_{T2})\,.$$

Above the helicity-dependent off-shell matrix elements were calculated as:

$$\begin{aligned} M(\lambda_{W^+}\lambda_{W^-}) &= \frac{1}{|\vec{q}_{\perp 1}||\vec{q}_{\perp 2}|} \sum_{\lambda_1 \lambda_2} (\vec{e}_\perp(\lambda_1) \cdot \vec{q}_{\perp 1})(\vec{e}_\perp^{\,*}(\lambda_2) \cdot \vec{q}_{\perp 2}) \mathcal{M}(\lambda_1, \lambda_2; \lambda_{W^+}, \lambda_{W^-}) \\ &= \frac{1}{|\vec{q}_{\perp 1}||\vec{q}_{\perp 2}|} \sum_{\lambda_1 \lambda_2} q^i_{\perp 1} q^j_{\perp 2}\, e_i(\lambda_1) e^*_j(\lambda_2) \mathcal{M}(\lambda_1, \lambda_2; \lambda_{W^+}, \lambda_{W^-})\,. \quad (2) \end{aligned}$$

Initial- and final-state helicity-dependent matrix elements were discussed e.g. in [11]. The k_t-factorization W-boson helicity dependent matrix elements were calculated with the help of the above [8].

The unintegrated inelastic flux of photons is expressed as:

$$\begin{aligned} \mathcal{F}^{\rm in}_{\gamma^* \leftarrow A}(z, \vec{q}_T) &= \frac{\alpha_{\rm em}}{\pi} \Bigg\{ (1-z) \left[\frac{\vec{q}_T^{\,2}}{\vec{q}_T^{\,2} + z(M_X^2 - m_p^2) + z^2 m_p^2} \right]^2 \frac{F_2(x_{\rm Bj}, Q^2)}{Q^2 + M_X^2 - m_p^2} \\ &+ \frac{z^2}{4 x_{\rm Bj}^2} \frac{\vec{q}_T^{\,2}}{\vec{q}_T^{\,2} + z(M_X^2 - m_p^2) + z^2 m_p^2} \frac{2 x_{\rm Bj} F_1(x_{\rm Bj}, Q^2)}{Q^2 + M_X^2 - m_p^2} \Bigg\}\,, \quad (3) \end{aligned}$$

The main ingredients of the formula are the F_1 and F_2 proton structure functions.

contribution	8 TeV	13 TeV
LUX-like		
$\gamma_{el}\gamma_{in}$	0.214	0.409
$\gamma_{in}\gamma_{el}$	0.214	0.409
$\gamma_{in}\gamma_{in}$	0.478	1.090
ALLM97 F2		
$\gamma_{el}\gamma_{in}$	0.197	0.318
$\gamma_{in}\gamma_{el}$	0.197	0.318
$\gamma_{in}\gamma_{in}$	0.289	0.701
SU F2		
$\gamma_{el}\gamma_{in}$	0.192	0.420
$\gamma_{in}\gamma_{el}$	0.192	0.420
$\gamma_{in}\gamma_{in}$	0.396	0.927
LUXqed collinear		
$\gamma_{in+el}\ \gamma_{in+el}$	0.366	0.778
MRST04 QED collinear		
$\gamma_{el}\gamma_{in}$	0.171	0.341
$\gamma_{in}\gamma_{el}$	0.171	0.341
$\gamma_{in}\gamma_{in}$	0.548	0.980
Elastic- Elastic		
$\gamma_{el}\gamma_{el}$ (Budnev)	0.130	0.273
$\gamma_{el}\gamma_{el}$ (DZ)	0.124	0.267

Table 1: *Cross sections (in pb) for different contributions and different F_2 structure functions: LUX, ALLM97 and SU, compared to the relevant collinear distributions with MRST04 QED and LUXqed distributions.*

The unintegrated elastic flux of photons is expressed as:

$$\begin{aligned}\mathcal{F}^{\mathrm{el}}_{\gamma^*\leftarrow A}(z,\vec{q}_T) &= \frac{\alpha_{\mathrm{em}}}{\pi}\Bigg\{(1-z)\left[\frac{\vec{q}_T^{\,2}}{\vec{q}_T^{\,2}+z(M_X^2-m_p^2)+z^2m_p^2}\right]^2\frac{4m_p^2G_E^2(Q^2)+Q^2G_M^2(Q^2)}{4m_p^2+Q^2}\\ &+ \frac{z^2}{4}\frac{\vec{q}_T^{\,2}}{\vec{q}_T^{\,2}+z(M_X^2-m_p^2)+z^2m_p^2}G_M^2(Q^2)\Bigg\}\ .\end{aligned} \tag{4}$$

In this case the main ingredients are G_E and G_M, the proton electromagnetic form factors.

To calculate the inelastic fluxes of photons, one needs the numerical representation of the proton structure functions. Different parametrizations of the F_2 structure functions are available in the literature, see e.g. [12, 13, 14].

3 Results

The integrated cross sections obtained in our approach are collected in Table 1.

Without any gap survival effects one has:

$$\sigma(inel.-inel.) > \sigma(inel.-el.) + \sigma(el.-inel.) > \sigma(el.-el.)\ . \tag{5}$$

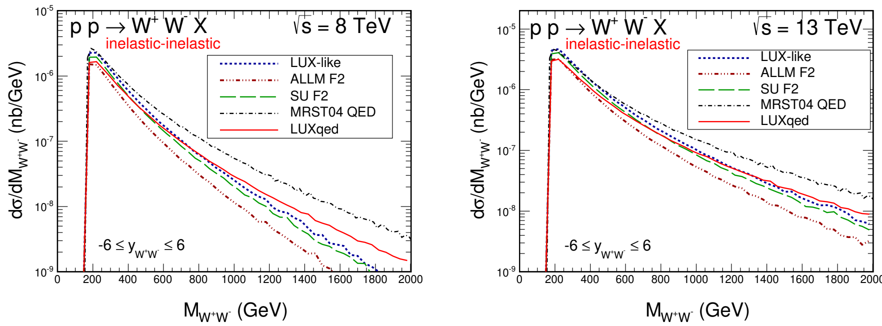

Figure 2: *M_{WW} invariant mass distribution for double dissociative contribution obtained with different parametrizations of structure functions.*

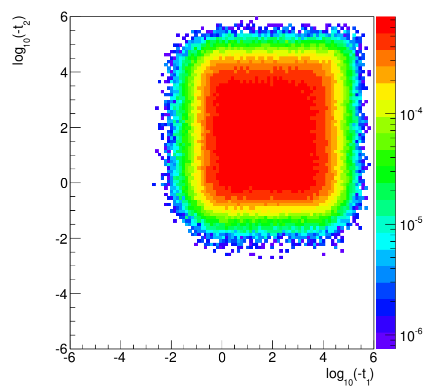

Figure 3: *Two-dimensional distribution in* $(\log_{10}(Q_1^2), \log_{10}(Q_2^2))$ *for double-dissociative processes.*

Many differential distributions were calculated in [8]. Here, in Fig. 3, we show only the invariant-mass distribution for double-dissociation processes (inelastic-inelastic) for different parametrizations of the structure functions taken from the literature.

The k_t-factorization result is similar to the collinear one for the same structure function (LUX-like). The rather old MRST04-QED collinear approach [15] predicted larger cross section. The reasons were discussed in [8].

As an example, in Fig. 3 we show the distribution of the photon virtualities. Rather large photon virtualities come into the game. Such large virtualities seem to contradict collinear approach.

The remnant fragmentation [9] was done with the help of PYTHIA 8 program. Including only parton (jet) emission is already a quite good approximation.

The gap survival probability for single dissociative process is calculated as:

$$S_R(\eta_{\rm cut}) = 1 - \frac{1}{\sigma}\int_{-\eta_{\rm cut}}^{\eta_{\rm cut}} \frac{{\rm d}\sigma}{{\rm d}\eta_{\rm jet}} {\rm d}\eta_{\rm jet} \,. \tag{6}$$

Jet emissions were considered also in [17]; the gap survival factor associated with jet emission is shown in Fig. 4. We find (see also Table 1)

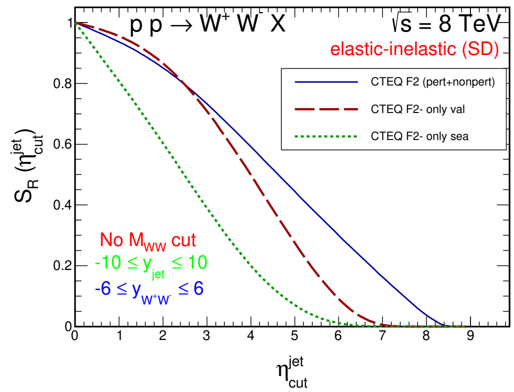

Figure 4: *Gap survival factor for single dissociative process associated with the jet emission. The solid line is for the full model, the dashed line for the valence contribution and the dotted line for the sea contribution.*

$$S_{R,DD} \approx (S_{R,SD})^2 \ . \tag{7}$$

Such an effect is expected when the two fragmentations are independent, which is the case in the model construction. So far we have not included the soft gap survival factors. They are relatively easy to calculate only for double-elastic (DE) contribution [16]. For the "soft" gap survival factors we expect:

$$S_{soft}(DD) < S_{soft}(SD) < S_{soft}(DE) \ . \tag{8}$$

	8 TeV	13 TeV	8 TeV	13 TeV	8 TeV	13 TeV
$(2M_{WW}, 200$ GeV$)$	0.763(2)	0.769(2)	0.582(4)	0.591(4)	0.586(1)	0.601(2)
$(200, 500$ GeV$)$	0.787(1)	0.799(1)	0.619(2)	0.638(2)	0.629(1)	0.649(1)
$(500, 1000$ GeV$)$	0.812(2)	0.831(2)	0.659(3)	0.691(3)	0.673(2)	0.705(2)
$(1000, 2000$ GeV$)$	0.838(7)	0.873(5)	0.702(12)	0.762(8)	0.697(5)	0.763(6)
full range	0.782(1)	0.799(1)	0.611(2)	0.638(2)	0.617(1)	0.646(1)

Table 2: *Average rapidity gap survival factors:* $S_{R,SD}(|\eta^{ch}| < 2.5)$, $(S_{R,SD})^2\,(|\eta^{ch}| < 2.5)$, $S_{R,DD}(|\eta^{ch}| < 2.5)$ *related to remnant fragmentation for single dissociative and double dissociative contributions for different ranges of* M_{WW}.

Finally we wish to show also similar results for the $pp \to t\bar{t}$ reaction. In Table 3 we show the integrated cross section for different categories of processes. Rather small cross sections are obtained; it is not clear at present whether such a process can be identified experimentally.

As an example we show $t\bar{t}$ invariant mass distribution for inclusive case as well as when extra veto on (mini)jet is imposed. The inclusion of rapidity gap veto reduces the cross section. Whether the cross section corresponding to the photon-photon fusion can be measured requires special dedicated studies.

4 Conclusions

Helicity-dependent matrix elements for $\gamma^*\gamma^* \to W^+W^-$ (off-shell photons) have been derived and used in the calculation of cross sections for $pp \to W^+W^-$ reaction. We have obtained a cross section of about 1 pb at the LHC energies. Different combinations of the final states (elastic-elastic, elastic-inelastic, inelastic-elastic, inelastic-inelastic) have been considered. Several correlation observables have been studied. Large

Contribution	No cuts	$y_{\rm jet}$ cut
elastic-elastic	0.292	0.292
elastic-inelastic	0.544	0.439
inelastic-elastic	0.544	0.439
inelastic-inelastic	0.983	0.622
all contributions	2.36	1.79

Table 3: *Cross section for $t\bar{t}$ production in fb at $\sqrt{s}$ = 13 TeV for different components, without (left column) and with the extra condition on the outgoing jet $|y_{\rm jet}| > 2.5$ (right).*

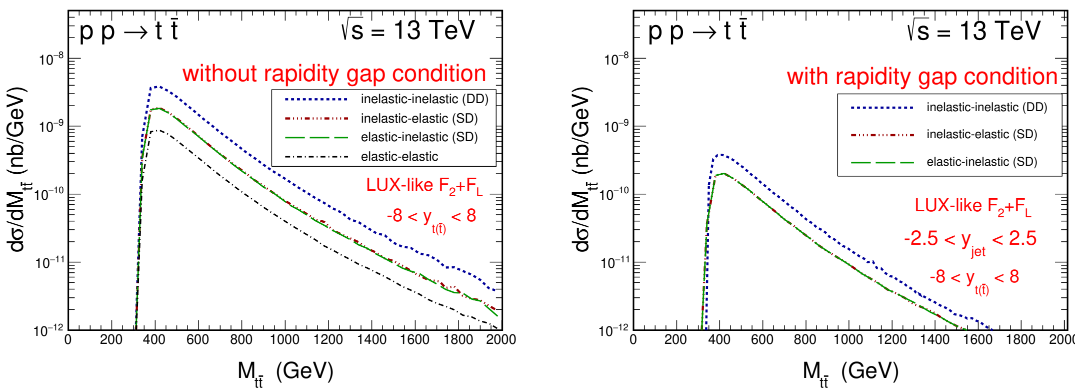

Figure 5: *$t\bar{t}$ invariant-mass distribution for different components defined in the figure. The left panel is without imposing the condition on the struck quark/antiquark and the right panel includes the condition.*

contributions from the regions of large photon virtualities Q_1^2 and/or Q_2^2 have been found putting in question the reliability of leading-order collinear-factorization approach.

We have discussed the quantity called "remnant gap survival factor" for the $pp \to W^+W^-$ reaction initiated via photon-photon fusion. We have calculated the gap survival factor for single dissociative process at parton level. In such an approach the outgoing parton (jet/mini-jet) is responsible for destroying the rapidity gap. We have found that the hadronisation only mildly modifies the gap survival factor calculated at parton level. We have found different values for double- and single-dissociative processes. In general, $S_{R,DD} < S_{R,SD}$ and $S_{R,DD} \approx (S_{R,SD})^2$. We expect that the factorisation observed here for the remnant dissociation and hadronisation will be violated when the soft processes are explicitly included. The larger $\eta_{\rm cut}$ (upper limit on charged particles pseudorapidity), the smaller rapidity gap survival factor S_R. This holds for both double and single dissociation. The present approach is a first step towards a realistic modelling of gap survival in photon induced interactions and definitely requires further detailed studies and comparisons to the existing and future experimental data. We have shown that rather large photon virtualities come into the game for W^+W^- production.

We have also calculated the cross sections for $t\bar{t}$ production via $\gamma\gamma$ mechanism in pp collisions, including the photon transverse momenta and using modern parametrizations of the proton structure functions. The contribution to the inclusive $t\bar{t}$ is only about 2.5 fb. We have found $\sigma_{tt}^{ela-ela} < \sigma_{tt}^{SD} < \sigma_{tt}^{DD}$. We have calculated several differential distributions. Some of them are not accessible in the standard equivalent-photon approximation. As for W^+W^- production, we have shown that rather large photon virtualities come into the game.

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