## THE ERGODICITY OF THE MOTION OF A PARTICLE IN A POTENTIAL FIELD

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## 1. Introduction

The ergodicity of classical dynamical systems which appear actually in the statistical mechanics was discussed by Ya. G. Sinai [11]. He announced the ergodicity of the dynamical system of particles with central potentials of a special type which are enclosed in a rectangular box. However no proofs have been published yet, except the simplest one-particle model which is called a Sinai billiard system. The ergodicity of the system and the K-property are shown in [13], [8], and the Bernoulli property is shown in [5]. The purpose of this report is to show the ergodicity and the Bernoulli property of the motion of a particle in a potential field. The proofs for a special case of a compound central potential field are separately presented in [9], [10].

## 2. Observation

Let T be a two dimensional torus and let  $U_1(q)$ , i = 1, 2, ..., I, be several potential functions with finite ranges  $\overline{Q}_1$ 's defined on T. Suppose that the potential ranges do not overlap and that the boundary  $\partial Q_1$  of the range  $\overline{Q}_1$  is a closed curve of  $C^3$ -class. Assume that every  $U_1(q)$  is continuous in the torus T and is continuously differentiable in  $\overline{Q}_1$ .

Observe the motion of a particle with mass m and energy E in the potential field. Denote by  $\{S_t\}$  the flow induced from the dynamical system; that is, for each (q, p) the point  $S_t(q, p)$  means the state of the particle at time t whose initial state is (q, p), where  $q = (q^{(1)}, q^{(2)})$  and  $p = (p^{(1)}, p^{(2)})$  are the position and the momentum. As usual the flow  $\{S_t\}$  can be restricted to the energy surface  $M_E$ ,  $M_E \equiv \{(q, p) ; (p^{(1)})^2 + (p^{(2)})^2 = 2m(E-U(q)), q \in Q_E\}$ , where  $Q_E \equiv \{q ; U(q) \le E\}$ . Moreover the measure

$$d\mu_{\rm F} \equiv {\rm const.} d\omega dq^{(1)} dq^{(2)}$$

on  $M_E$  is invariant under  $\{S_t\}$ , where  $(p^{(1)}, p^{(2)}) = (\kappa \cos \omega, \kappa \sin \omega)$ with  $\kappa^2 = 2m(E-U(q))$ . Let  $\pi$  be the natural projection from  $M_E$  to I the configuration space  $Q_E$ ;  $\pi(q, p) = q$ . Put  $Q \equiv T - \bigcup_{\substack{i=1 \\ i=1}}^{u} Q_i$  and  $M_0 \equiv \pi^{-1}(Q)$ . It is easily seen that almost every orbit of the particle crosses the boundary  $\partial Q = \bigcup \partial Q_1$  of Q. A point q of the boundary can be parametrized by  $(\iota, r)$ , where  $\iota$  is the number of the curve  $\partial Q_1$ which contains q, and r is the arclength between the point q and a fixed origin of  $\partial Q_1$  measured along the curve  $\partial Q_1$  clockwise. Further, a point (q, p) in the set  $\pi^{-1}(\partial Q)$  can be parametrized by  $(\iota, r, \varphi)$ , where  $\varphi$  is the angle between p and the inward normal of  $\partial Q_1$  at q. Put

$$M \equiv \{(\iota, r, \boldsymbol{\varphi}) \in \pi^{-1}(\partial Q) ; \frac{1}{2}\pi \leq \boldsymbol{\varphi} \leq \frac{3}{2}\pi\},\$$

namely M is the set of all incident vectors at  $\partial Q$ . Then a transformation  $T_*$  of M is defined by that  $T_*^{-1}(\iota, r, \mathcal{G}) = (\iota_1, r_1, \mathcal{G}_1)$ shows the next incident position and angle after starting from  $(\iota, r, \mathcal{G})$ . Then it is seen that  $T_*$  preserves the measure  $\nu$  of M defined by

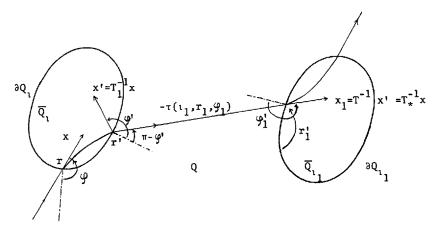
$$dv = -v_0 \cos g d g dr du$$

with normalized constant  $v_0$ , where  $d_1$  means the unit mass on each  $\iota$ .

Since it is true that  $\{S_t\}$  is ergodic if and only if  $T_*$  is ergodic, it is important to investigate the transformation  $T_*$ . Then the transformation  $T_*$  can be resolved into the product

$$\mathbf{T}_* = \mathbf{T}_1 \mathbf{T},$$

where  $T_1$  is the transformation such that for  $(\iota, r', g') \equiv T_1^{-1}(\iota, r, g)$ , the point  $(\iota, r', \pi - g')$  shows the position and the momentum at the instant when the particle goes out from  $\overline{Q}_1$ , and T is the Sinai billiard transformation which maps  $T_{\pi}^{-1}(\iota, r, \mathcal{G}) = (\iota_1, r_1, \mathcal{G}_1)$  to  $(\iota_1, r', \mathcal{G}')$ .



Generally, let us call a transformation  $T_{*}$  a perturbed billiard transformation if  $T_{*}$  is expressed in the form

 $T_* = T_T$ 

with the Sinai billiard transformation T defined for Q and with a  $C^2$ -diffeomorphism  $T_1$  of M which preserves the measure  $\nu$  and each set  $\pi^{-1}(\partial Q_1)$ ,  $\iota = 1, 2, \cdots, I$ .

Obviously, the transformation  $T_*$  induced from  $\{S_t\}$  is a perturbed billiard transformation. If the perturbation of T by  $T_1$  is small, then the ergodic theoretical properties of T should be preserved under the perturbation.

3. Theorems

In this section, assume the following assumptions (H-1)  $\sim$  (H-3): (H-1) Every  $\overline{\mathbb{Q}}_1$  is a strictly convex domain with the boundary  $\Im_{\mathbb{Q}}$  of  $\mathbb{C}^3$ -class.

(H-2)  $T_{\rm l}$  is a  ${\rm C^2-diffeomorphism}$  of M which preserves the measure  $\nu$  and satisfies

$$T_1(1, r, \frac{1}{2}\pi) = (1, r, \frac{1}{2}\pi)$$
 and  $T_1(1, r, \frac{3}{2}\pi) = (1, r, \frac{3}{2}\pi).$ 

Denote by  $k(\iota, r)$  the curvature of  $Q_{\iota}$  at r and by  $-\tau(\iota_1, r_1, \mathscr{G}_1)$  the distance between  $(\iota_1, r_1)$  and  $(\iota, r')$  measured along the orbit with  $(\iota, r', \mathscr{G}') = T(\iota_1, r_1, \mathscr{G}_1)$ . Put  $A \equiv \frac{\partial r_1}{\partial r}$ ,  $B \equiv \frac{\partial r_1}{\partial \mathscr{G}}$ ,  $C \equiv \frac{\partial \mathscr{G}_1}{\partial r}$  and  $D \equiv \frac{\partial \mathscr{G}_1}{\partial \mathscr{G}}$ .

(H-3) There exist positive constants n,  $L_{max}$ ,  $L_{min}^{+}(1)$  and  $L_{min}^{-}(1)$ ,  $1 \le i \le I$ , such that for  $(1, r, \mathcal{G})$  with  $(1, r', \mathcal{G}') = T_{1}^{-1}(1, r, \mathcal{G})$  and  $(1_{1}, r_{1}, \mathcal{G}_{1}) = T_{*}^{-1}(1, r, \mathcal{G})$  and for  $s \ge L_{min}^{+}(1)$  (resp.  $s \ge L_{min}^{-}(1)$ ) the following inequalities hold

$$-Cs-D \ge 1+\eta \qquad (resp. -\frac{\cos \mathscr{P}_1}{\cos \mathscr{P}} (Cs+A) \ge 1+\eta),$$

$$L_{max} \ge \frac{As+B}{Cs+D} \ge L_{min}^+(\iota_1) \qquad (resp. L_{max} \ge \frac{Ds+B}{Cs+A} \ge L_{min}^-(\iota)),$$

$$\frac{\cos \mathscr{P}_1}{\cos \mathscr{P}} (A+\frac{1}{s}B)(Cs+D) \ge 1+\eta \qquad (resp. \frac{\cos \mathscr{P}_1}{\cos \mathscr{P}} (D+\frac{1}{s}B)(Cs+A) \ge 1+\eta).$$

The assumption (H-3) may look rather complicated. However the assumption is satisfied if the perturbation by  $T_1$  is small. For example, the following proposition gives sufficient conditions under which the assumption (H-3) is satisfied. For a function f, put  $[f]^+ \equiv \max\{f, 0\}$ .

Proposition 1.

(i) If there exists a positive constant  $\eta$  such that

$$-D \ge 1+n$$
,  $-\frac{\cos g}{\cos g} A \ge 1+n$ 

and if  $-B \ge 0$ ,  $-C \ge 0$ , then the assumption (H-3) is fulfilled. (ii) If

$$\frac{1}{k_{\min}} \left[ -\frac{\partial g'}{\partial r} \right]^{+} + \left( 1 + \frac{1}{k_{\min} |\tau|_{\min}} \right) \left[ -\frac{\partial g'}{\partial g} + 1 \right]^{+} < 1,$$

$$k_{\max} \left( 1 + \frac{2}{k_{\min} |\tau|_{\min}} \right) \left[ \frac{\partial g}{\partial r} \right]^{+} + \left( 1 + \frac{1}{k_{\min} |\tau|_{\min}} \right) \left[ -\frac{\partial g'}{\partial g'} + 1 \right]^{+} < 1,$$

then the assumption (H-3) is fulfilled, where  $k_{\min} \equiv \min_{\iota, r} k(\iota, r)$ ,

 $\underset{i,r}{\operatorname{kmax}} = \underset{i,r}{\operatorname{max}} k(i, r) \text{ and } |\tau|_{\min} = \underset{i,r,\mathcal{G}}{\operatorname{min}} |\tau(i, r, \mathcal{G})|.$ 

Theorem 2. Under the assumptions (H-1), (H-2) and (H-3), the transformation  $T_* = T_1 T$  is Bernoullian, and hence it is ergodic and a K-system.

Theorem 3. Let  $\{S_t\}$  be the flow given in §1 which is governed by potential functions  $\{U_t\}$ . If  $\{U_t\}$  and the energy E of the particle admit the assumptions (H-1), (H-2) and (H-3), then  $\{S_t\}$  is ergodic. Moreover, if  $\{S_t\}$  has no point spectrum, then  $\{S_t\}$  is a K-system and a Bernoulli flow.

If  $U_{l}$ 's are central potentials, then one can state the sufficient condition for the ergodicity in terms of  $\{U_{l}\}$ . A central potential V(s) is called *bell-shaped* if

(BS-1) V(s) is continuous for s > 0 and V(s) = 0 for s > R with some R,

(BS-2) V(s) belongs to  $C^2$ -class in (0, R) and there exist left derivatives V'(R-0) and V"(R-0),

(BS-3) -sV'(s) is monotone decreasing and V'(R-0) < 0.

Theorem 4. If every  ${\rm U}_{\rm l}$  is bell-shaped and if the energy E satisfies the condition

$$0 < E < \frac{1}{4} \min_{1} \{-\frac{R_{1}|\tau|_{\min}}{R_{1}+|\tau|_{\min}} U'(R_{1}-0)\},\$$

then  $\{S_t\}$  is ergodic. Moreover if  $\{S_t\}$  has no point spectrum, then  $\{S_t\}$  is a K-system and a Bernoulli flow.

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