

TRANSVERSE MOMENTUM OF DILEPTONS

NA3 Collaboration

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ABSTRACT

Transverse momentum produced by π^- at 150, 200 and 280 GeV are compared with QCD predictions.

RESUME

Les impulsions transverses de dimuons produits par des π^- de 150, 200 et 280 GeV, sont comparées aux prédictions de CDQ.

Incident π^- of 150 GeV, 200 GeV and 280 GeV producing muon pairs on a platinum target were studied in the NA3 spectrometer. The distribution of the transverse momentum p_t is expected to be described by the convolution of two functions : the first one is the distribution of the quark intrinsic momentum k_t ; the second is deduced from QCD calculations. In terms of mean squared values, one has :

$$\langle p_t^2 \rangle = \langle k_t^2 \rangle + \langle p_t^2 \rangle_{\text{QCD}}$$

Without scaling violations, one expects :

$$\langle p_t^2 \rangle_{\text{QCD}} = S f(\tau, y), \quad \tau = M^2/S$$

y is the rapidity and M the mass of the dimuon.

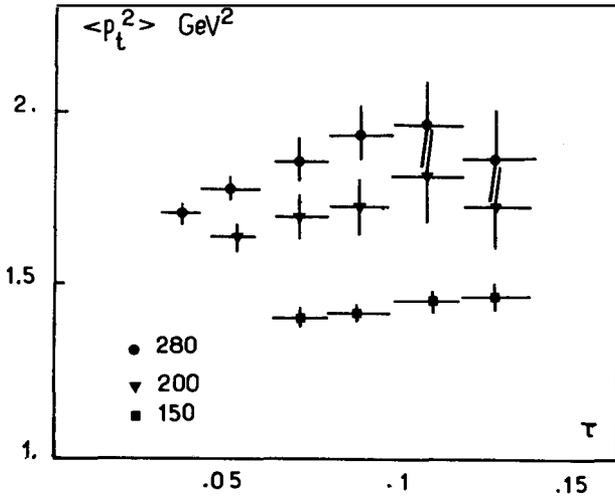


Fig. 1 : $\langle p_t^2 \rangle$ versus $\tau = M^2/S$

The τ dependance of $\langle p_t^2 \rangle$ is presented in figure 1, after integration over y in the range $[-.3, 1.1]$. The quoted errors are statistical only, but systematic effects of the order of $.1 \text{ GeV}^2$ have to be added. They are essentially induced by the beam spread in the target. The data are consistent with the expected form :

$$\langle p_t^2 \rangle = .85 + S F (\tau)$$

In figure 2, one has plotted the quantity : $(\langle p_t^2 \rangle - .85)/S$. A comparison with QCD calculations may be performed. One calls σ_{DY} the cross section corresponding to the graph :

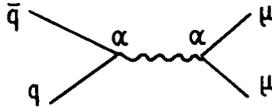
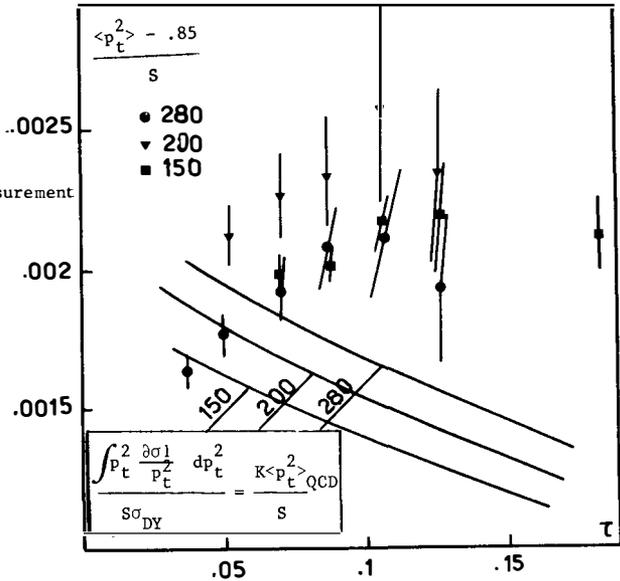
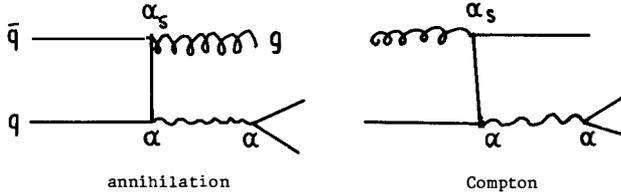


Fig. 2
Experimental measurement compared to QCD calculations.



It is possible to calculate $\frac{d\sigma_1}{2 dp_t}$ corresponding to the first order in α_s . It involves external gluons as in the graphs :



In π^- interactions, the annihilation term is predominant contrary to the proton case. The total cross section is related to σ_{DY} by the K factor²⁾

$$\sigma_{\text{tot}} = K \sigma_{DY}$$

Then, at the first order in α_s :

$$\langle p_t^2 \rangle_{\text{QCD}} = \frac{\int_{p_t^2}^2 \frac{d\sigma_1}{dp_t^2} dp_t^2}{K \sigma_{DY}}$$

In order to take into account scaling violation, one has chosen :

$$\alpha_s = 12 \pi / 25 \text{ Log } (1 + 4 Q^2)$$

The quarks structure functions are deduced from deep inelastic³⁾ and Drell-Yan^{4,5)} experiments :

Proton :

$$\begin{aligned} u &= 2A (\alpha, \beta) x^\alpha (1-x)^\beta + \bar{u} \\ d &= A(\alpha, \beta+1) x^\alpha (1-x)^{\beta+1} + \bar{d} \\ \bar{u} &= \bar{d} = 2\lambda = 2\bar{\lambda} = A(1-x)^n \\ \alpha &= .52 - .16 \bar{s} \\ \beta &= 2.79 + .77 \bar{s} \\ A &= .26 + .18 \bar{s} \\ n &= 7.80 - .78 \bar{s} \end{aligned}$$

$$\bar{s} = \text{Log } \frac{\text{Log } 4 Q^2}{\text{Log } 80}$$

Pion :

$$\begin{aligned} \bar{u} &= d = A(\alpha', \beta') x^{\alpha'} + u \\ u &= \bar{d} = 2\lambda = 2\bar{\lambda} = A'(1-x)^{n'} \\ \alpha' &= .5 - .1 \bar{s} \\ \beta' &= 1. + .7 \bar{s} \\ A' &= .12 + .7 \bar{s} \\ n' &= 5. \end{aligned}$$

$$A(\alpha, \beta) = \frac{\Gamma(\alpha + \beta + 1)}{\Gamma(\alpha) \Gamma(\beta + 1)}$$

The smaller contribution due to the compton-graphs is computed with gluon structure functions of the form : $Ag(1-x)^{ng}$, with :
 $Ag = 3.49 + 4.1 \bar{S}$, $ng = 5.75 \pm 7.5 \bar{S}$ for the nucleon and $Ag = 2.$, $ng = 3.$ for the pion.

The results of the calculation are displayed in figure 2, where $\langle p_t^2 \rangle_{\text{QCD}}$ is multiply by the K factor. The confrontation with experimental values suggests that $\frac{d\sigma}{dp_t^2}$ has to be corrected by a factor greater than this K factor.. Furthermore, the QCD estimations decrease with τ , whereas experimental values increase. In conclusion, one has :

$$\langle p_t^2 \rangle = \langle k_t^2 \rangle \frac{\int_{p_t^2}^2 \frac{d\sigma}{dp_t^2} dp_t^2 + \int_{p_t^2}^2 \frac{d\sigma_n}{dp_t^2} dp_t^2}{K \sigma_{DY}}$$

with $\langle k_t^2 \rangle = .85 \pm .1$

$$\frac{d\sigma_n}{dp_t^2} > K \frac{d\sigma}{dp_t^2}$$

$\frac{d\sigma_n}{dp_t^2}$ represents all the contributions which have to be added to the first order of QCD.

In figure 3, $\langle p_t^2 \rangle$ is displayed as a function of y , after integration over τ , in the range [.06, .14]. Shapes agree with theoretical ones resulting essentially from phase space.

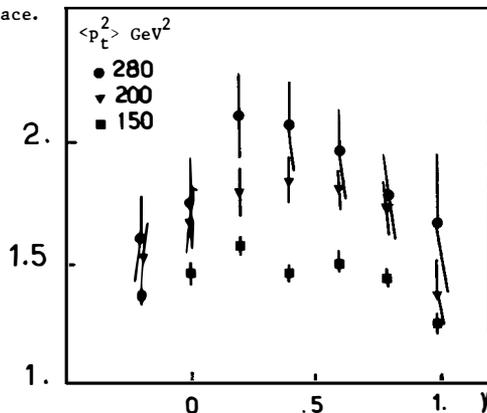


Figure 3
 $\langle p_t^2 \rangle$ versus center of mass rapidity

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