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The Noodle Method*

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ABSTRACT

The Noodle method allows a fast and efficient sampling of the region of integration during Monte Carlo calculations. The method works by automatically optimizing the weight distribution and by replacing analytic results by exact numerical calculations. The method allows integrations of certain functions that have been heretofore considered infeasible.

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1. Introduction

The traditional Monte Carlo method¹ is unacceptably inefficient for certain commonly occurring integrands and requires various analytic formulae in addition to the integrand. The Noodle method eliminates these complications. The remainder of this section contains a critical description of the traditional Monte Carlo method, section 2 describes the Noodle Method, section 3 contains a detailed example of an application of the Noodle method, namely a numerical computation of a physical cross-section, and section 4 compares the Noodle method to the traditional Monte Carlo method.

Consider the integral

$$I = \int_0^1 dx \int_0^1 dy \int_0^1 dz f(x, y, z) \quad (1)$$

The Monte Carlo method is based on the following equation:

$$\frac{Vol(S)}{N} \sum_{i=1}^{i=N} f(x_i) = \int_S f \pm \sqrt{\frac{Var(f)}{N}} \quad (2)$$

where each x_i is a randomly and uniformly chosen point in the space S .

Consequently, to estimate I accurately for small values of N , I is first rewritten as

$$\int_0^1 dx \int_0^1 dy \int_0^{G(x,y)} dG \frac{f(x, y, z(G))}{g(x, y, z(G))}, \text{ where } G(x, y) = \int_0^1 dz g(x, y, z) \quad (3)$$

with the corresponding Monte Carlo sum

$$\frac{Vol(\tilde{S})}{N} \sum_{i=1}^{i=N} \frac{f(x_i, y_i, z(G_i))}{g(x_i, y_i, z(G_i))} \pm \sqrt{\frac{Var(\frac{f}{g})}{N}} \quad (4)$$

where each (x_i, y_i, G_i) is chosen randomly and uniformly from \tilde{S} (figure 1). Then

g is chosen so that $\text{Var}(\frac{f}{g})$ is small.

For this method to work, g has to satisfy two mutually antagonistic criteria. It is crucial that g has a rather complicated form so that it approximates f well. Yet at the same time, g has to have a simple form so that it can be integrated analytically. Choosing such a g is an art form and is the complication with the traditional Monte Carlo method that impedes a person from applying the method.

Once g has been chosen, the Monte Carlo sum is performed by first generating a $(x_i, y_i, z(G_i))$ triplet, then by evaluating the summand. The usual algorithm for generating (x_i, y_i, G_i) s is to generate a random point inside the smallest box that contains \tilde{S} and reject the point if it lies outside \tilde{S} . This algorithm becomes inefficient when \tilde{S} has "spikes" (*figure 2*). Moreover, it is not easy to invert

$$G_i = \int_0^{z(G_i)} dz g(x, y, z) \quad (5)$$

to evaluate $z(G_i)$.

2. The Noodle Method

In the Noodle method, S is decomposed into small sets over each of which f is roughly constant. Then for each set, g is defined as constant and equals f at least one point in the set. One such choice of g for a simple decomposition of S is

$$g(x, y, z) = f\left(\frac{i}{M}, \frac{j}{M}, 0\right); \frac{i}{M} < x < \frac{i+1}{M}, \frac{j}{M} < y < \frac{j+1}{M}, 0 < z < 1 \quad (6)$$

for some M and all $0 \leq i, j < M$. The corresponding \tilde{S} is shown on *figure 3a*. The decomposition of S induces a decomposition of \tilde{S} via the change of

coordinates z into G :

$$\frac{i}{M} < x < \frac{i+1}{M}, \frac{j}{M} < y < \frac{j+1}{M}, 0 < G < G_{i,j}; G_{i,j} = \frac{1}{M^2} g\left(\frac{i}{M}, \frac{j}{M}\right).$$

Each of these sets is called a noodle. These noodles are then arbitrarily ordered from 1 up to the total number of noodles.

By construction, g approximates f , and, since g is constant over each noodle, g can be integrated exactly. This solves the problem of finding a suitable g .

Moreover, defining g locally constant allows a fast and efficient generation of the $(x_i, y_i, z(G_i))$ triplets. This is possible because of the simple relation that the uniform distribution in \tilde{S} is equivalent to the distribution g in S . Consider $\rho(0 < x - x_0 < \delta x, 0 < y - y_0 < \delta y, 0 < G - G_0 < \delta G) = \delta x \delta y \delta G$, which is the uniform distribution in \tilde{S} . Since $G - G_0 = g(x, y, z_0)(z - z_0)$, we have $\rho(0 < x - x_0 < \delta x, 0 < y - y_0 < \delta y, 0 < z - z_0 < \frac{\delta G}{g}) = \delta x \delta y \delta G$. Since δG is small but otherwise arbitrary, replacing δG by $g \delta z$ yields the desired result. Using this relationship, we directly generate (x_i, y_i, z_i) with the distribution g thus eliminating the need for complicated inversion of functions. First, we arrange the noodles from end to end to make a long noodle by defining an array as follows:

NOODLE(0)= 0

NOODLE(i)= NOODLE($i - 1$)+volume of the i^{th} noodle

the length of the long noodle = NOODLE(number of the noodles)

This is shown in *figure 3b*. Finally, in order to generate a (x_i, y_i, z_i) , generate a random number between 0 and the length of the long noodle. This picks out a noodle uniformly in \tilde{S} and a set according to the distribution g in S . Once a noodle is chosen, since g is constant within the set, any point within the set is

equally likely to be chosen. Accordingly, we simply pick a random point (x, y, z) within the set. It is fast because each generation of (x, y, z) involves only a binary search algorithm for choosing a box and three random number generations. It is efficient in the sense that no point that is generated is rejected.

For this method to work better than the traditional Monte Carlo method, it is crucial to partition S into small sets so that f is indeed approximately constant over each set. One way to achieve this is to choose M large enough so that, over S , $\frac{f}{g}$ takes values in a small interval around 1, say $(1 - \epsilon, 1 + \epsilon)$. Consequently, in our example, $\frac{1}{M}$ is the maximum scale in $x - y$ space over which the fractional change in f is less than ϵ . Thus, if the graph of f has a tall but narrow spike at a point but is otherwise flat, the necessary value of M dictated by the width of the spike is clearly an overestimate for the rest of the region of integration.

In the spirit of replacing nontrivial analytical calculations by exact numerical calculations, it would be pleasing if there existed a numerical algorithm for finding an acceptable partition of S . The following algorithm is only a partial solution to this problem that presumes a knowledge of the general profile of the integrand. The completely general numerical solution is usually unnecessary and seems to be very complicated. It should be mentioned that there are integration routines, VEGAS for example, that solve this problem in general, but their generality is usually compensated by their slow speed.

```
procedure bakery;  
    index:= 1;  
    while (index ≤ length of the noodle) do  
        begin
```

```

while is_wild( $index^{th}$  noodle) do
begin
    partition the noodle into smaller noodles;
    let one of the noodles be the  $index^{th}$  noodle;
    put the rest of the new noodles at the end of the ordering;
    increase the number of noodles by the number of new noodles;
end;
noodle(index):= noodle(index-1)+ volume of the  $index^{th}$  noodle;
index:= index+1;
end;
end;

function is_wild(noodle):boolean;
is_wild:=( $f$  fluctuates unacceptably inside the box);
end;

```

The user has to define the *wild* function. Because the Noodle method automatically samples S , a list of points in S that are associated with unacceptable weights can readily be generated. This furnishes two pieces of information: Unacceptable weights indicate that the *wild* function is insufficiently defined, and the list indicate the regions of S over which the integrand has unexpected structures. Using this information, the *wild* function can be redefined so that the unacceptable weights do not occur.

3. An Example

We illustrate the method by using the noodle method to compute the $e^+e^- \rightarrow \mu^+\mu^-$ cross section after the effects of the initial state bremmstrahlung have been taken into account. The cross section is given by²

$$\sigma = \int_{x_{cut}}^1 dx_+ \int_{x_{cut}}^1 dx_- \int d\Omega D(x_+)D(x_-) \frac{d\sigma_0}{d\Omega_{CM}}(\theta_{cm}, x_+x_-S) \Theta(p_{\mu^+}^{lab}, p_{\mu^-}^{lab}) \quad (7)$$

where σ_0 is the uncorrected cross-section, and $\Theta(x) = 1$ if x is inside the phase space cut and 0, otherwise. For the sake of simplicity, we consider the case of no phase space cut.

In order to integrate this numerically, it is necessary to regulate the poles of D by a suitable change of variables. The resulting integral is

$$\int_0^{P_{cut}} dP_+ \int_0^{P_{cut}} dP_- \int_0^{2\pi} d\phi \int_{-1}^1 d\cos\theta w_d(P_+)w_d(P_-) \frac{d\sigma}{d\Omega}(\cos\theta, x_+(P_+)x_-(P_-)S) \quad (8)$$

where $P_+ = (1 - x_+)^{\frac{1}{2}}$, the singular part of the integrated electron structure function. We shall evaluate this integral for $\sqrt{S} = 94\text{Gev}$. The corresponding \tilde{S} is shown on *figure 4a*. Since the D s are always approximately 1, the only *wild* regions are when $\sqrt{x_+x_-S}$ is near the *Z-pole* or the *photon-pole*. Moreover, the noodles should not be too big. Hence, it is natural to cut S evenly in the x_+x_- -plane and to consider the two extreme values of x_-x_+S to decide whether a particular noodle is wild.

Following the Noodle method, at the beginning of the program, we make noodles:

```

procedure bakery;
arrays  $x_{-}^{min}$ ,  $x_{-}^{max}$ ,  $x_{+}^{min}$ ,  $x_{+}^{max}$ , NOODLE
number of noodles:= 1;
index:= 1;
 $x_{-}^{min}$ (index):=  $x_{cut}$ ;
 $x_{+}^{min}$ (index):=  $x_{cut}$ ;
 $x_{-}^{max}$ (index):= 1;
 $x_{+}^{max}$ (index):= 1;
while (index ≤ number of the noodles) do
  while is_wild( $x_{-}^{min}$ (index) $x_{+}^{min}$ (index)S,  $x_{-}^{max}$ (index) $x_{+}^{max}$ (index)S) do
    begin
       $x_{-}^{min}$ (number of noodles+1):=  $x_{-}^{min}$ (index);
       $x_{+}^{max}$ (number of noodles+1):=  $x_{+}^{max}$ (index);
       $x_{-}^{max}$ (number of noodles+1):=  $\frac{1}{2}(x_{-}^{max}(index)+x_{-}^{min}(index))$ ;
       $x_{+}^{min}$ (number of noodles+1):=  $\frac{1}{2}(x_{+}^{max}(index)+x_{+}^{min}(index))$ ;
       $x_{-}^{min}$ (number of noodles+2):=  $x_{-}^{max}$ (number of noodles+1);
       $x_{+}^{max}$ (number of noodles+2):=  $x_{+}^{max}$ (index);
       $x_{-}^{max}$ (number of noodles+2):=  $x_{-}^{max}$ (index);
       $x_{+}^{min}$ (number of noodles+2):=  $x_{+}^{min}$ (number of noodles+1);
       $x_{-}^{min}$ (number of noodles+3):=  $x_{-}^{min}$ (number of noodles+2);
       $x_{+}^{max}$ (number of noodles+3):=  $x_{+}^{min}$ (number of noodles+2);
       $x_{-}^{max}$ (number of noodles+3):=  $x_{-}^{max}$ (number of noodles+2);
       $x_{+}^{min}$ (number of noodles+3):=  $x_{+}^{max}$ (number of noodles+2);
       $x_{+}^{max}$ (index):=  $x_{+}^{max}$ (number of noodles+3);
       $x_{-}^{max}$ (index):=  $x_{-}^{max}$ (number of noodles+1);
    end;
  end;
end;

```

```

    number of noodles:= number of noodles + 3;

end;

 $x_+ := \frac{1}{2}(x_{min}^{min} + x_{max}^+);$ 
 $x_- := \frac{1}{2}(x_{min}^{min} + x_{max}^-);$ 
 $P_+ := (1 - x_+)^{\frac{\beta}{2}};$ 
 $P_- := (1 - x_-)^{\frac{\beta}{2}};$ 
 $\Delta P_+ := (1 - x_{min}^+)^{\frac{\beta}{2}} - (1 - x_{max}^+)^{\frac{\beta}{2}};$ 
 $\Delta P_- := (1 - x_{min}^-)^{\frac{\beta}{2}} - (1 - x_{max}^-)^{\frac{\beta}{2}};$ 
NOODLE(index):=
    NOODLE(index-1)+ $w_d(P_-)w_d(P_+)\frac{d\sigma}{d\Omega}(\cos\theta, x_+x_-S)\Delta P_+ \Delta P_- \Delta\cos\theta;$ 
end;

```

```

function is_wild( $S_{min}, S_{max}$ );
    wild if  $S_{min} < M_Z^2 + 9\Gamma_Z M_Z$  and  $M_Z^2 - 9\Gamma_Z M_Z < S_{max}$ ;
    wild if  $S_{min} < S_\gamma$ ;
    wild if  $\sqrt{S_{max}} - \sqrt{S_{min}} \geq \Delta E$ ;
    otherwise not wild;
end;

```

The sampling of the integration region proceeds in two steps. First, we pick a noodle with the relative probability of NOODLE(index). Second, we pick a random point $(P_i^+, P_i^-, \cos\theta_i)$ within the noodle:

```

procedure generate weights;
    pick a random number  $r$  between 0 and the length of the noodle;
    find  $l$  such that  $\text{NOODLE}(l-1) < r \leq \text{NOODLE}(l)$ ;
 $\Delta P^+ := (1 - x_{min}^+(l))^{\frac{\beta}{2}} - (1 - x_{max}^+(l))^{\frac{\beta}{2}};$ 

```

```

 $\Delta P^- := (1 - x_{min}^-(l))^{\frac{\beta}{2}} - (1 - x_{max}^-(l))^{\frac{\beta}{2}};$ 
 $r_1$  and  $r_2$  are random numbers between 0 and 1;
 $P_+ := (1 - x_{max}^+)^{\frac{\beta}{2}} + r_1 \Delta P^+;$ 
 $P_- := (1 - x_{max}^-)^{\frac{\beta}{2}} + r_2 \Delta P^-;$ 
 $x_+ := 1 - P_+^{\frac{2}{\beta}};$ 
 $x_- := 1 - P_-^{\frac{2}{\beta}};$ 
weight:=  $\frac{w_d(P_-)w_d(P_+) \frac{d\sigma}{dt}(\cos\theta_l, x_+ x_- S) \Delta P^+ \Delta P^- \Delta \cos\theta}{NOODLE(l) - NOODLE(l-1)};$ 
end;

```

4. Implementation and Comparison

The simplicity and the efficiency of the Noodle method makes it trivial to write programs to integrate quite complicated functions and to investigate the properties of the integral in a short period of time. An implementation of the preceding algorithm appears in the *Appendix*. For 10^4 Monte Carlo samples, we reproduce the results obtained by Berends et al. at the Z -peak.³ The resulting weight distribution is shown on *figure 5*. The results of a more complete calculation that contains the initial state radiation, the weak vertex functions, and the vector boson self-energies appears elsewhere⁴.

As discussed earlier, the Noodle method is a simplification of the Monte Carlo method in two aspects. First, it is no longer necessary to supply the approximate integrands to perform integrations using the Monte Carlo method. Second, there is no need for complicated inversion formulae. More importantly, however, the Noodle method is an improvement over the Monte Carlo method in that it can perform certain integrations that are not feasible to do by the traditional Monte Carlo method. In computations of e^+e^- to f^+f^- at energies around 1Tev, the

shape of S becomes quite singular as shown on *figures 4b and 4c*. The ratio of the volume of the smallest box that contains S to the volume of S is approximately 10^4 and in order to find the value of σ_{lab} to 1 percent, the number of samples required in the traditional Monte Carlo method is 10^8 , compared to only 10^4 in the Noodle method. The associated weight distribution is shown on *figure 6*.

5. Acknowledgement

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REFERENCES

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2. O. Nicrosini and L. Trentadue, Phys. Lett. 196B:551, 1987
3. Berends, Burgers, van Neerman, Phys. Lett. 185B:395,1987
4. Kennedy, Lynn, Stuart, Im, SLAC-PUB-4128

6. Captions

Figure 1

The change of variables from z to G improve the convergence rate of the Monte Carlo sum.

Figure 2

The conventional Monte Carlo Method becomes inefficient for certain shapes of the space of integration, such as the one shown here.

Figure 3

In the Noodle method, the integrand is approximated by step functions. The resulting \tilde{S} , corresponding to the same integrand as in figure 1, is shown on the left. The blocks are then arranged into the shape of a noodle, as shown on the right.

Figure 4

The graphs of \tilde{S} at 94GeV (a), 110GeV (b), 210GeV (c) for integrand appearing in equation 8 are shown here. Each group of blocks of the same height, the group of blocks appearing at the $(P_+, P_-) = (1, 1)$ corner on figure 4a, for example, corresponds to a single noodle. This correspondence is explicitly illustrated on figure 4c. In order to show detail, only the portion of $P_+ P_-$ in which both coordinates are greater than 0.8 is shown, and the photon pole, appearing at $(P_+, P_-) = (1, 1)$ corner and the Z^0 pole, the other peak in each graph, are truncated.

Figure 5

The distribution of the ratio between the integrand in equation 8 and the approximate integrand constructed using noodles at $\sqrt{S} = 94\text{GeV}$.

Figure 6

The distribution of the ratio between the integrand in equation 8 and the approximate integrand constructed using noodles at $\sqrt{S} = 1TeV$.

7. Appendix

```

C-----C
C THE NOODLE MONTE CARLO - EXPONENTIATED INITIAL STATE BREMS- C
C STRAHLUNG FOR THE Z0 RESONANCE C
C E+E- STRUCTURE FUNCTIONS BY O. NICROSINI & L. TRENTADUE C
C MONTE CARLO METHOD AND CODING BY J. IM C
C REVISIONS AND COMMENTS BY D. KENNEDY C
C SLAC SLC/MARK II - 10/30/87 C
C NEW IMPROVED VERSION - 2/19/88 C
C
C
C CONTAINS SUBROUTINES THAT CUT NOODLES ALONG GIVEN XP AND XM C
C START WITH ARBITRARY PARTITION OF THE X-SPACE C
C CUT NOODLES ALONG XP AND XM LINES AT UNIFORM INTERVALS C
C HAND IT OVER TO WILD FUNCTION C
C CONTAINS THETA CUT C
C-----C
C-----C
C IMPLICIT REAL*8 (A-Z)
C INTEGER I,NEVENT,NN,MM,STAT,FILE,BIGRAD
C DIMENSION MUP(4),MUM(4),CUT(5)
C COMMON/COUPL/CEL,CER,CML,CMR,QE,QM
C COMMON/PARAM/BETA,ALPHA,PI,ME,MZ,GAMZ,SINW2,CONST,DELP,CUT,POL,
C SCUT,SLPOL,SHPOL
C-----PARAMETERS
C
C: MONTE CARLO PARAMETERS: NUMBER OF EVENTS, NUMBER OF NOODLES = NN^2
C: NEVENT= 20000
C: NN= 5
C: MM= 5
C: CALL RAN11A(37561547)
C: PHYSICS CONSTANTS
C: PI= 3.141592D0
C: ALPHA= 1.D0/137.03602D0/(1.D0-.06D0)
C: SINW2= 0.23D0
C: CONST= ALPHA**2/4.D0*389385.7D0
C: ME= .000511D0
C: MZ= 92.D0
C: GAMZ= 2.9D0
C: ELECTRON AND MUON COUPLINGS
C: QE= -1.D0
C: QM= -1.D0
C: I3E= -.5D0
C: I3M= -.5D0
C: CEL= 2.D0*(I3E-SINW2*QE)
C: CER= 2.D0*(-SINW2*QE)
C: CML= 2.D0*(I3M-SINW2*QM)
C: CMR= 2.D0*(-SINW2*QM)
C: BEAM ENERGY AND PHASE SPACE CUTS
C: ECUT IS A COMMON ENERGY CUT THAT SHOULD BE EQUAL TO OR LESS THAN
C: THE SMALLER OF THE TWO CUT(1) AND CUT(2)
C: EB2= 92.D0
C: DEB2= 0.5D0
C: MAXEB2=91.D0
C: ECUT= 5.D0
C: SCUT= (45.D0)**2
C: SHPOL= 7.D0*MZ*GAMZ+MZ**2
C: SLPOL= -7.D0*MZ*GAMZ+MZ**2

```

```

CUT(1)= 42.4D0
CUT(2)= 42.4D0
CUT(3)= 1.D0
CUT(4)= DCOS(PI*6.07D0/180.D0)
CUT(5)= 1.D0
STAT= 1
FILE= 1
C-----MAIN
C: LEFT POL, THEN RIGHT POL
POL= 1.D0
23 S= EB2**2
C: INITIALIZES AND CONSTRUCTS NOODLE FOR A GIVEN S
CALL INIT(S)
RADCUT= DSQRT(S*(1.D0-(BETA/(1.D0+BETA))))
24 POL= -1.D0*POL
CALL BAKERY(S,NN,MM,ECUT,APRCRS)
MAXW= 0.D0
FBASYM= 0.D0
LRASYM= 0.D0
WSUM= 0.D0
WFSUM= 0.D0
WBSUM= 0.D0
WERR= 0.D0
WFERR= 0.D0
WBERR= 0.D0
SIGERR= 0.D0
BIGRAD= 0
C: GENERATES NEVENT EVENTS
DO 100 I= 1,NEVENT
IF (I-10000*(I/10000) .EQ. 0) PRINT 101,I
101 FORMAT(' ',I10,' EVENTS PROCESSED')
CALL GENEV(S,SP,MUP,MUM,W,DFLOAT(I)/NEVENT)
IF (DSQRT(SP) .LT. RADCUT) THEN
    BIGRAD= BIGRAD+1
ENDIF
IF (POL .LT. 0.D0) CALL HISTO(W,1.D0,0.D0,8.D0,1)
IF (POL .GT. 0.D0) CALL HISTO(W,1.D0,0.D0,8.D0,2)
WSUM= WSUM+W
SIGERR= SIGERR+W**2
C: BINS WEIGHTS IN WEIGHT DISTRIBUTION HISTOGRAMS AND CHECKS FOR
C: MAXIMUM WEIGHTS
IF ((STAT .EQ. 1).AND.(W .GT. MAXW)) MAXW= W
IF (MUM(3).GT.0.D0) THEN
    WFSUM= WFSUM+W
    WFERR= WFERR+W**2
ELSE
    WBSUM= WBSUM+W
    WBERR= WBERR+W**2
ENDIF
100 CONTINUE
C: MONTE CARLO INTEGRATION VOLUME, APR AND EXACT CROSS SECTIONS
DV= APRCRS/NEVENT
SIG= WSUM*D
ERR= DV*DSQRT(SIGERR-WSUM**2/NEVENT)
C: FORWARD/BACKWARD ASYMMETRY
FBASYM= (WFSUM-WBSUM)/(WFSUM+WBSUM)
FBAERR= WBSUM**2*WFERR+WFSUM**2*WBERR

```

```

FBAERR= DSQRT(FBAERR)/(WFSUM+WBSUM)**2
IF (POL .LT. 0.D0) THEN
  PRINT 1,EB2,SIG,ERR
  PRINT 2,FBASYM,FBAERR
  IF (STAT .EQ. 1) PRINT 3,WSUM/NEVENT,ERR/DV/NEVENT,MAXW,BIGRAD,
    RADCUT
  IF (FILE .EQ. 1) THEN
    WRITE (24,9) EB2,SIG,ERR
    WRITE (27,9) EB2,FBASYM,FBAERR
  ENDIF
  APPL= APRCRS
  SIGL= SIG
  ERRL= ERR
  WFSUML= WFSUM*APPL
  WBSUML= WBSUM*APPL
  WFERRL= WFERR*APPL**2
  WBERRL= WBERR*APPL**2
  GOTO 24
ELSE
  PRINT 4,EB2,SIG,ERR
  PRINT 5,FBASYM,FBAERR
  IF (STAT .EQ. 1) PRINT 6,WSUM/NEVENT,ERR/DV/NEVENT,MAXW,BIGRAD,
    RADCUT
  APPR= APRCRS
  SIGR= SIG
  ERRR= ERR
  WFSUMR= WFSUM*APPR
  WBSUMR= WBSUM*APPR
  WFERRR= WFERR*APPR**2
  WBERRR= WBERR*APPR**2
  SIGTOT= 0.5D0*(SIGL+SIGR)
  ERRTOT= 0.5D0*DSQRT(ERRL**2+ERRR**2)
  LRASYM= (SIGL-SIGR)/(SIGL+SIGR)
  LRAERR= DSQRT((SIGR*ERRL)**2+(SIGL*ERRR)**2)/2.D0/SIGTOT**2
  FBATOT= (WFSUML+WFSUMR-WBSUML-WBSUMR)/
    (WFSUML+WFSUMR+WBSUML+WBSUMR)
  FBERRT= (WBSUML+WBSUMR)**2*(WFERRL+WFERRR)
  FBERRT= FBERRT+(WFSUML+WFSUMR)**2*(WBERRL+WBERRR)
  FBERRT= DSQRT(FBERRT)/(WFSUML+WFUMR+WBSUML+WBSUMR)**2
  PRINT 7,EB2,SIGTOT,ERRTOT,FBATOT,FBERRT
  PRINT 8,LRASYM,LRAERR
  IF (FILE. EQ. 1) THEN
    WRITE (25,9) EB2,SIGR,ERRR
    WRITE (28,9) EB2,FBASYM,FBAERR
    WRITE (23,9) EB2,SIGTOT,ERRTOT
    WRITE (26,9) EB2,FBATOT,FBERRT
    WRITE (29,9) EB2,LRASYM,LRAERR
  ENDIF
  EB2= EB2+DEB2
ENDIF
C: OUTPUTS CROSS SECTION, L/R AND F/B ASYMMETRIES, ERRORS
1 FORMAT(' LEFT: ROOT(S)=',F9.4,'      SIGL: ',F10.8,'+-',F6.4)
2 FORMAT(' FBASYML: ',F10.8,'+-',F6.4)
3 FORMAT(' STATISTICS: SUM(WL)=',F10.4,'    SUM(WL2)=',F10.4,
      '    MAX(WL)=',F10.4,'/',
      '    I6, ' EVENTS BELOW ',F9.4,' GEV', '/')
4 FORMAT(' RIGHT: ROOT(S)=',F9.4,'      SIGR: ',F10.8,'+-',F6.4)

```

```

5 FORMAT(' FBASYMR: ',F10.8,'+-',F6.4)
6 FORMAT(' STATISTICS: SUM(WR)=' ,F10.4,'    SUM(WR2)=' ,F10.4,
        .      MAX(WR)=' ,F10.4,/,
        .      ',I6,' EVENTS BELOW ',F9.4,' GEV',/)
7 FORMAT(' TOTAL: ROOT(S)=' ,F9.4,'          SIG: ',F10.8,'+-',F6.4,/,
        .      FBASYM: ',F10.8,'+-',F6.4,/)
8 FORMAT(' LRASYM: ',F10.8,'+-',F6.4,/)
9 FORMAT(F10.4,'   ',F14.8,'   ',F14.8)
IF (EB2 .LE. MAXEB2) GOTO 23
IF ((STAT .EQ. 1).AND.(FILE .EQ. 1)) THEN
    CALL REPORT(3)
ENDIF
CALL REPORT(1)
CALL REPORT(2)
STOP
END

```

C-----INIT

```

SUBROUTINE INIT(S)
IMPLICIT REAL*8 (A-Z)
INTEGER I,J
DIMENSION CUT(5)
COMMON/PARAM/BETA,ALPHA,PI,ME,MZ,GAMZ,SINW2,CONST,DELP,CUT,POL,
SCUT,SLPOL,SHPOL

```

C: DEFINES PARAMETERS IN TRENTADUE'S FORM FACTORS

```

BETA= 2.D0*ALPHA/PI*(DLOG(S/ME**2)-1.D0)
Z2= 1.6449340668D0
Z3= 1.202056903159D0
L= DLOG(S/ME**2)
DELP2= 1.D0+BETA**2*PI**2/24.D0+ALPHA/PI
     *(1.5D0*L+PI**2/3.D0-2.D0)
     +ALPHA**2.D0/PI**2.D0*((9.D0/8.D0-2.D0*Z2)*L**2.D0
     +(-45.D0/16.D0+
     11.D0/2.D0*Z2+3.D0*Z3)*L-6.D0/5.D0*Z2**2.D0-9.D0/2.D0*Z3-
     6.D0*Z2*DLOG(2.D0)+3.D0/8.D0*Z2+57.D0/12.D0)
DELP= DSQRT(DELP2)

```

C: SETS UP WEIGHT HISTOGRAM

```

CALL CLRHST
RETURN
END

```

C-----BAKERY

```

SUBROUTINE BAKERY(S,NN,MM,ECUT,APRCRS)
IMPLICIT REAL*8 (A-Z)
INTEGER A,B,I,J,K,L,DL,NN,MM,PWRO2,THETA,LENGTH
DIMENSION NOODLE(50000),XP(50000),XM(50000),SIZP(50000),
         THETA(50000),SIZM(50000),CUT(5)
COMMON/PARAM/BETA,ALPHA,PI,ME,MZ,GAMZ,SINW2,CONST,DELP,CUT,POL,
SCUT,SLPOL,SHPOL
COMMON/BAKRY/NOODLE,XP,XM,SIZP,SIZM,THETA,L

```

C: CONSTRUCTS NOODLES - DISCRETE SPACE: APPRSIGO, APRP+, APRP-
C: NOODLES EXIST ONLY IN BAKERY AND ENTRY POINT GENWT BELOW

```

XCUT=ECUT*2.D0/DSQRT(S)
PCUT= GETP(XCUT)
L= 0
C: THE BIG SQUIRE
DO 1 I= 0,MM-1
    L= L+1
1
```

```

XP(L)= XCUT
XM(L)= XCUT
SIZP(L)= 1.D0-XCUT
SIZM(L)= 1.D0-XCUT
1 THETA(L)= I
C: THE LEFT RECTANGLE
DO 2 I= 0,MM-1
L= L+1
XP(L)= XCUT**2
XM(L)= XCUT
SIZP(L)= XCUT
SIZM(L)= (1.D0-XP(L))-XCUT
2 THETA(L)= I
C: THE RIGHT RECTANGLE
DO 3 I= 0,MM-1
L= L+1
XM(L)= XCUT**2
XP(L)= XCUT
SIZM(L)= XCUT
SIZP(L)= (1.D0-XM(L))-XCUT
3 THETA(L)= I
PRINT *,L
DO 321 I= 1,NN
CALL CUTXP(GETX(PCUT*I/NN))
321 CALL CUTXM(GETX(PCUT*I/NN))
LASTND= 0.D0
I= 1
23 NXP= XP(I)+SIZP(I)
NXM= XM(I)+SIZM(I)
C: FINDS WILD NOODLES - PHOTON AND Z PEAKS
WD= WILD(S*XP(I)*XM(I),S*NXP*NXM)
IF ((WD.GT.0.D0).AND.(L.LE.49996)) THEN
XP(L+1)= (XP(I)+NXP)/2.D0
XM(L+1)= XM(I)
THETA(L+1)= THETA(I)
XP(L+2)= XP(L+1)
XM(L+2)= (XM(I)+NXM)/2.D0
THETA(L+2)= THETA(I)
XP(L+3)= XP(I)
XM(L+3)= XM(L+2)
THETA(L+3)= THETA(I)
SIZP(I)= SIZP(I)/2.D0
SIZM(I)= SIZM(I)/2.D0
SIZP(L+1)= SIZP(I)
SIZM(L+1)= SIZM(I)
SIZP(L+2)= SIZP(I)
SIZM(L+2)= SIZM(I)
SIZP(L+3)= SIZP(I)
SIZM(L+3)= SIZM(I)
L= L+3
GOTO 23
ENDIF
24 XPMID= XP(I)+SIZP(I)/2.D0
XMMID= XM(I)+SIZM(I)/2.D0
CALL GETWD(XPMID,W1)
CALL GETWD(XMMID,W2)
CS= 2.D0*(THETA(I)+.5D0)/MM-1.D0

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DV= (GETP(XP(I)+SIZP(I))-GETP(XP(I)))*
     (GETP(XM(I)+SIZM(I))-GETP(XM(I)))*2.D0/MM
C: VECTOR NOODLE DEFINED AS CUMULATIVE SUM OF APPROXIMANT CROSS SECTIONS
NOODLE(I)= LASTND+DSIG(S*XPMID*XMMID,CS,POL)*W1*W2*D
LASTND= NOODLE(I)
I= I+1
IF (I.LE.L) GOTO 23
PWRO2= IDINT(DLOG(L+.5D0)/DLOG(2.D0))
LENGTH= L
PRINT *, LENGTH
C: FINAL ENTRY OF NOODLE = TOTAL CUMULATIVE MEGANOODLE
      - APR TOT CS/2PI
APRCRS= NOODLE(L)*2.D0*PI
RETURN
C-----GENWGT
ENTRY GENWGT(S,P,M,C,WGT,INDEX)

C: GENERATES EVENTS BY PICKING RANDOM POINT ALONG MEGA NOODLE
RN= INDEX*NOODLE(LENGTH)
C: FINDS CORRESPONDING NOODLE THRU BINARY SEARCH
L= 2**PWRO2
IF (L.GT.LENGTH) L= L/2
DL= L/2
111 IF ((RN.GT.NOODLE(L)).AND.(L+DL.LE.LENGTH)) THEN
    L= L+DL
ELSE IF (RN.LT.NOODLE(L-1)) THEN
    L= L-DL
ELSE IF (L+DL.LE.LENGTH) THEN
    GOTO 400
ENDIF
DL= DL/2
IF (DL.GT.0) GOTO 111
400 SIZEP= GETP(XP(L))-GETP(XP(L)+SIZP(L))
SIZEM= GETP(XM(L))-GETP(XM(L)+SIZM(L))
C: GENERATES EVENT: X+, X-, ANGLE, WEIGHT
P= GETX(GETP(XP(L))-RANDOM(SIZEP))
M= GETX(GETP(XM(L))-RANDOM(SIZEM))
CALL GETWD(P,WP)
CALL GETWD(M,WM)
C= 2.D0*THETA(L)/MM-1.D0+RANDOM(2.D0/MM)
WGT= WP*WM*DSIG(S*P*M,C,POL)*SIZEP*SIZEM*2.D0/MM/
      (NOODLE(L)-NOODLE(L-1))
C: BINS S/PRIME DISTRIBUTION
RETURN
END
C-----CUTXP
SUBROUTINE CUTXP(X)
IMPLICIT REAL*8 (A-Z)
INTEGER I,L,THETA
DIMENSION NOODLE(50000),XP(50000),XM(50000),SIZP(50000),
          THETA(50000),SIZM(50000),CUT(5)
COMMON/BAKRY/NOODLE,XP,XM,SIZP,SIZM,THETA,L

DO 124 I=1,L
  IF ((XP(I).LT.X).AND.(XP(I)+SIZP(I).GT.X)) THEN
    L= L+1
    XP(L)= X

```

```

        XM(L)= XM(I)
        SIZP(L)= XP(I)+SIZP(I)-X
        SIZM(L)= SIZM(I)
        THETA(L)= THETA(I)
        SIZP(I)= X-XP(I)
    ENDIF
124 CONTINUE
RETURN
END
C-----CUTXM
SUBROUTINE CUTXM(Y)
IMPLICIT REAL*8 (A-Z)
INTEGER I,L,THETA
DIMENSION NOODLE(50000),XP(50000),XM(50000),SIZP(50000),
           THETA(50000),SIZM(50000),CUT(5)
COMMON/BAKRY/NOODLE,XP,XM,SIZP,SIZM,THETA,L

DO 124 I=1,L
    IF ((XM(I).LT.Y).AND.(XM(I)+SIZM(I).GT.Y)) THEN
        L= L+1
        XM(L)= Y
        XP(L)= XP(I)
        SIZM(L)= XM(I)+SIZM(I)-Y
        SIZP(L)= SIZP(I)
        THETA(L)= THETA(I)
        SIZM(I)= Y-XM(I)
    ENDIF
124 CONTINUE
RETURN
END
C-----GETP
FUNCTION GETP(X)
IMPLICIT REAL*8 (A-Z)
COMMON/PARAM/BETA,ALPHA,PI,ME,MZ,GAMZ,SINW2,CONST,DELP,CUT,POL,
          SCUT,SLPOL,SHPOL

C: COMPUTES FUNCTION P GIVEN X
IF (1.D0-X .LT. 0.2D0**2.D0/BETA) THEN
    GETP= 0.D0
ELSE
    GETP= (1.D0-X)**(BETA/2.D0)
ENDIF
RETURN
END
C-----WILD
FUNCTION WILD(SL,SH)
IMPLICIT REAL*8 (A-Z)
DIMENSION CUT(5)
COMMON/PARAM/BETA,ALPHA,PI,ME,MZ,GAMZ,SINW2,CONST,DELP,CUT,POL,
          SCUT,SLPOL,SHPOL

C: DEFINES WILD NOODLES AT PHOTON AND Z PEAKS
WILD= -1.D0
DE= DSQRT(SH)-DSQRT(SL)
IF (((SL.LT.SCUT).AND.(DE.GT.4.D0)).OR.
    ((SL.LE.SHPOL).AND.(SH.GE.SLPOL).AND.(DE.GT.GAMZ/2.D0)).OR.
    (DE.GT.50.D0)) WILD= 1.D0

```

```

        RETURN
        END
C-----RANDOM
        FUNCTION RANDOM(X)
        REAL*8 X
        REAL*4 RAN11

        RANDOM= X*RAN11(0)
        RETURN
        END
C-----GENEV
        SUBROUTINE GENEV(S,SPRIME,MUP,MUM,WGT,INDEX)
        IMPLICIT REAL*8 (A-Z)
        DIMENSION MUP(4),MUM(4)
        DIMENSION CUT(5)
        COMMON/PARAM/BETA,ALPHA,PI,ME,MZ,GAMZ,SINW2,CONST,DELP,CUT,POL,
        SCUT,SLPOL,SHPOL

C: GENERATES EVENT: X+, X-, APR DIFF CS
        CALL GENWGT(S,XP,XM,COSTHP,WGT,INDEX)
        IF (WGT .EQ. 0.D0) THEN
          W= 0.D0
        ELSE
C: DEFINES C.M. (PRIMED) FRAME KINEMATICS: RANDOMLY DISTRIBUTED
C: POLAR AND AZIMUTHAL ANGLES
          SPRIME= S*XP*XM
          SINTHP= DSQRT(1.D0-COSTHP**2)
          COSPHP= DCOS(RANDOM(2.D0*PI))
          SINPHP= DSQRT(1.D0-COSPHP**2)
          BOOST= (XM-XP)/(XM+XP)
          GAMMA= 1.D0/DSQRT(1.D0-BOOST**2)
          EP= DSQRT(SPRIME)/2.D0
C: FIRST DEFINES MUM AND MUP MOMENTA IN C.M. (PRIMED) FRAME
          MUM(1)= EP*SINTHP*COSPHP
          MUM(2)= EP*SINTHP*SINPHP
          MUM(3)= EP*COSTHP
          MUM(4)= EP
          MUP(1)= -MUM(1)
          MUP(2)= -MUM(2)
          MUP(3)= -MUM(3)
          MUP(4)= MUM(4)
C: BOOSTS BACK TO LAB FRAME
          MP3= GAMMA*(MUP(3)+BOOST*MUP(4))
          MP4= GAMMA*(MUP(4)+BOOST*MUP(3))
          MM3= GAMMA*(MUM(3)+BOOST*MUM(4))
          MM4= GAMMA*(MUM(4)+BOOST*MUM(3))
C: THEN REDEFINES MUM AND MUP MOMENTA IN LAB (UNPRIMED) FRAME
          MUP(3)= MP3
          MUM(3)= MM3
          MUP(4)= MP4
          MUM(4)= MM4
C: DEFINES TOTAL LEFT AND RIGHT WEIGHTS WITH PHASE SPACE CUTS
          WGT= WGT*PHCUT(MUM,MUP)
          END IF
          RETURN
          END
C-----GETX

```

```

FUNCTION GETX(P)
IMPLICIT REAL*8 (A-Z)
DIMENSION CUT(5)
COMMON/PARAM/BETA,ALPHA,PI,ME,MZ,GAMZ,SINW2,CONST,DELP,CUT,POL,
SCUT,SLPOL,SHPOL
C: COMPUTES X GIVEN FUNCTION P
  IF (P .EQ. 1.D0) THEN
    GETX= 0.D0
  ELSE IF (P .LT. 0.2D0) THEN
    GETX= 1.D0
  ELSE
    GETX= 1.D0-P**2.D0/BETA
  END IF
  RETURN
END

C-----GETWD
SUBROUTINE GETWD(X,W)
IMPLICIT REAL*8 (A-Z)
DIMENSION CUT(5)
COMMON/PARAM/BETA,ALPHA,PI,ME,MZ,GAMZ,SINW2,CONST,DELP,CUT,POL,
SCUT,SLPOL,SHPOL

C: COMPUTES BREMSSTRAHLUNG WEIGHT = EXACT D/APR D
C: IF X = 0, THROW EVENT WAY
  IF (X .EQ. 0.D0) THEN
    W= 0.D0
  C: IF P < 0.1, X ESSENTIALLY ONE
    ELSE IF (X .GT. .9999999D0) THEN
      W= DELP
  C: OTHERWISE...
    ELSE
      W= DELP-.5D0*(1.D0+X)*(1.D0-X)**(1.D0-BETA/2.D0)+BETA/16.D0*
        ((X+1.D0)*(-4.D0*DLOG(1.D0-X)+3.D0*DLOG(X))-4.D0/(1.D0-X)
        *DLOG(X)-5.D0-X)*(1.D0-X)**(1.D0-BETA/2.D0)
    END IF
    RETURN
  END

C-----PHCUT
FUNCTION PHCUT(MUP,MUM)
IMPLICIT REAL*8 (A-Z)
DIMENSION MUP(4),MUM(4),CUT(5)
COMMON/PARAM/BETA,ALPHA,PI,ME,MZ,GAMZ,SINW2,CONST,DELP,CUT,POL,
SCUT,SLPOL,SHPOL

  IF (CUT(5).EQ.0.D0) THEN
    PHCUT= 1.D0
    RETURN
  ENDIF
C: MIN ENERGY CUT
  ECUT= 1.D0
  IF ((MUM(4).LT.CUT(1)).AND.(MUP(4).GE.CUT(2))) ECUT=0.D0
C: ENDCAP ANGLE CUT
  COSTH= MUM(3)/MUM(4)
  THCUT= 1.D0
  IF (DABS(COSTH).GT.CUT(3)) THCUT= 0.D0
C: ACOLLINARITY CUT
  DOT= MUM(1)*MUP(1)+MUM(2)*MUP(2)+MUM(3)*MUP(3)

```

```

COSXI= -DOT/MUM(4)/MUP(4)
XICUT= 1.D0
IF (COSXI.LT.CUT(4)) XICUT= 0.D0
C: TOTAL CUT
    PHCUT= ECUT*THCUT*XICUT
    RETURN
    END
C-----CLRHST
SUBROUTINE CLRHST
INTEGER I,J,N
REAL*8 HST,MIN,MAX,X,W,XMIN,XMAX
DIMENSION HST(500,5),MIN(5),MAX(5)

C: CLEARS COMMON HISTOGRAM
    DO 1 I= 1,500
    DO 1 J= 1,5
1 HST(I,J)= 0.D0
    RETURN
C-----HISTO
ENTRY HISTO(X,W,XMIN,XMAX,N)

C: BINS WEIGHTS IN HISTOGRAM
    MIN(N)= XMIN
    MAX(N)= XMAX
    I= (X-XMIN)/(XMAX-XMIN)*500
    IF ((I .LE. 500) .AND. (I .GT. 0)) HST(I,N)= HST(I,N)+W
    RETURN
C-----REPORT
ENTRY REPORT(N)

C: OUTPUTS HISTOGRAMS
    DO 2 I=1,500
2 WRITE (30+N,*) I*(MAX(N)-MIN(N))/500.D0+MIN(N),HST(I,N)
    RETURN
    END
C-----DSIG
FUNCTION DSIG(SP,X,POL)
IMPLICIT REAL*8 (A-Z)

C: DEFINES CROSS SECTION DEPENDING ON POLARIZATION
    IF (POL .EQ. -1.D0) DSIG= DSIGL(SP,X)
    IF (POL .EQ. 1.D0) DSIG= DSIGR(SP,X)
    RETURN
    END
C-----DSIGL
FUNCTION DSIGL(SP,X)
IMPLICIT REAL*8 (A-Z)
COMPLEX*16 R
DIMENSION CUT(5)
COMMON/COUPL/CEL,CER,CML,CMR,QE,QM
COMMON/PARAM/BETA,ALPHA,PI,ME,MZ,GAMZ,SINW2,CONST,DELP,CUT,POL,
SCUT,SLPOL,SHPOL

C: LEFT-HANDED CROSS SECTION - X = COS OF POLAR ANGLE
    DSIGL=CONST/SP*((1.D0+X)**2*(CDABS(QE*QM+R(SP)*CEL*CML))**2.D0
                  +(1.D0-X)**2*(CDABS(QE*QM+R(SP)*CEL*CMR))**2.D0)
    DSIGL=DSIGL/2.D0

```

```

        RETURN
        END
C-----DSIGR
        FUNCTION DSIGR(SP,X)
        IMPLICIT REAL*8 (A-Z)
        COMPLEX*16 R
        DIMENSION CUT(5)
        COMMON/COUPL/CEL,CER,CML,CMR,QE,QM
        COMMON/PARAM/BETA,ALPHA,PI,ME,MZ,GAMZ,SINW2,CONST,DELP,CUT,POL,
               SCUT,SLPOL,SHPOL

C: RIGHT-HANDED CROSS SECTION - X = COS OF POLAR ANGLE
        DSIGR=CONST/SP*((1.D0-X)**2*(CDABS(QE*QM+R(SP)*CER*CML))**2.D0
                      +(1.D0+X)**2*(CDABS(QE*QM+R(SP)*CER*CMR))**2.D0)
        DSIGR=DSIGR/2.D0
        RETURN
        END
C-----R
        FUNCTION R(S)
        IMPLICIT REAL*8 (A-Z)
        COMPLEX*16 R
        DIMENSION CUT(5)
        COMMON/PARAM/BETA,ALPHA,PI,ME,MZ,GAMZ,SINW2,CONST,DELP,CUT,POL,
               SCUT,SLPOL,SHPOL

C: NORMALIZED Z0 RESONANCE FACTOR
        R= S/4.D0/SINW2/(1.D0-SINW2)/DCMPLX(S-MZ**2,S*GAMZ/MZ)
        RETURN
        END

```

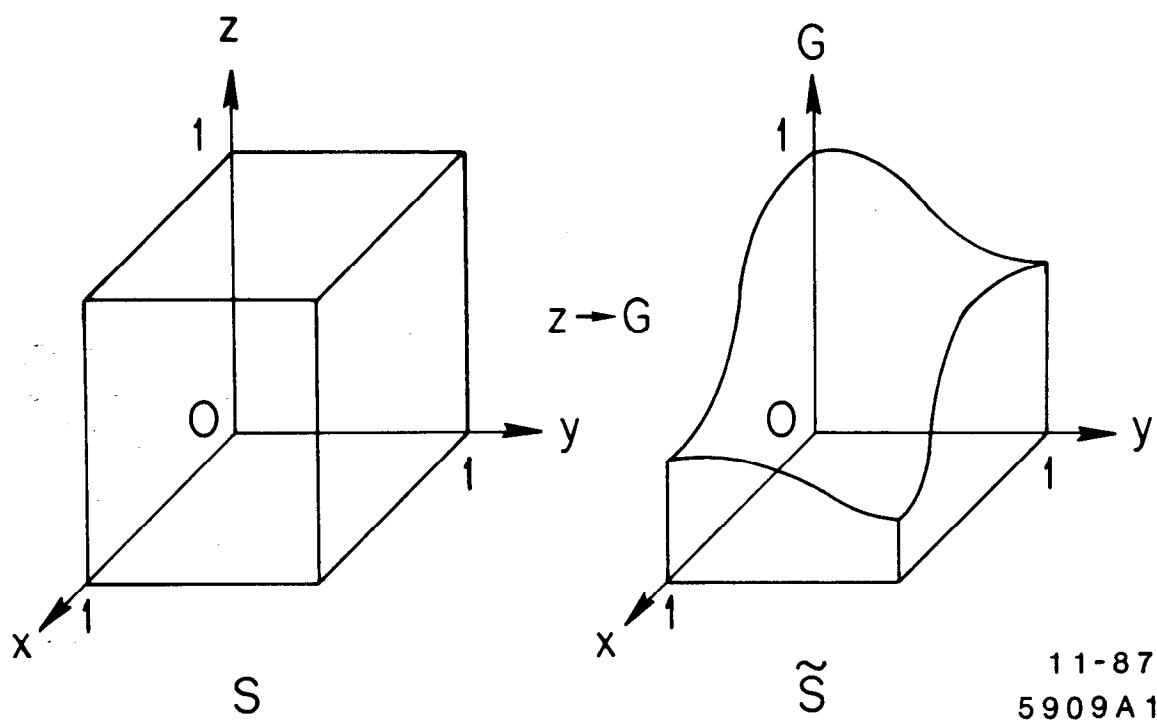
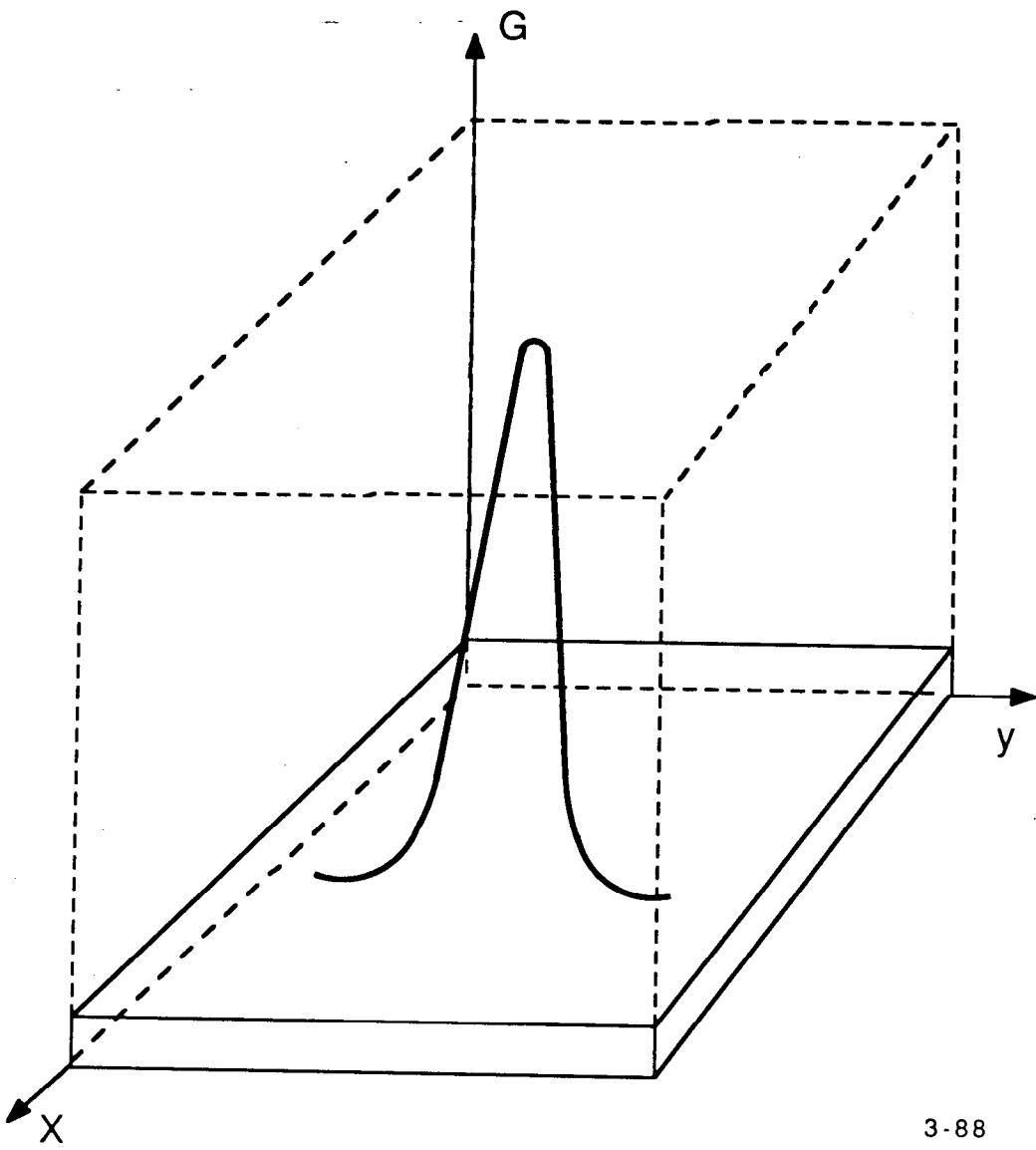
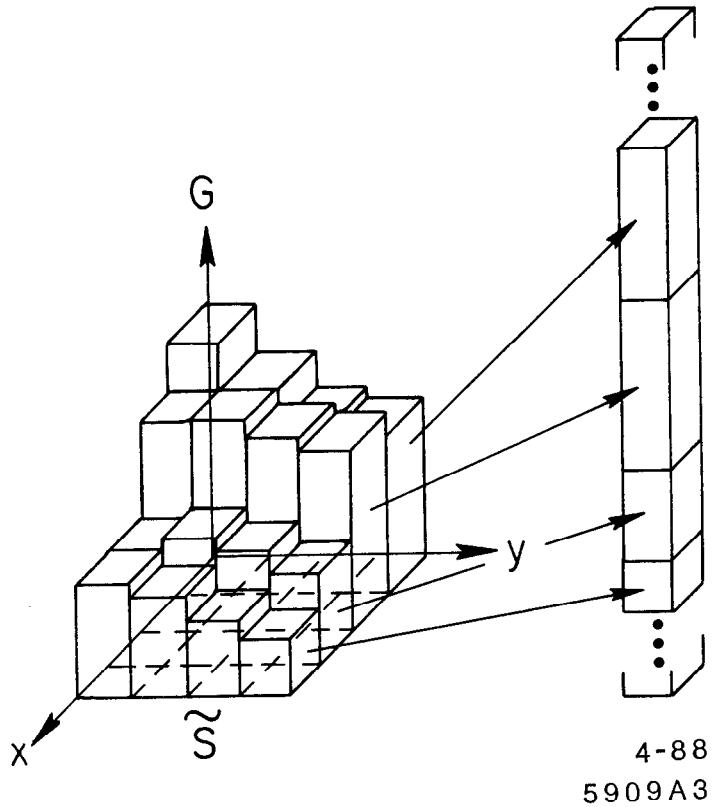


Fig. 1



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Fig. 2



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5909A3

Fig. 3

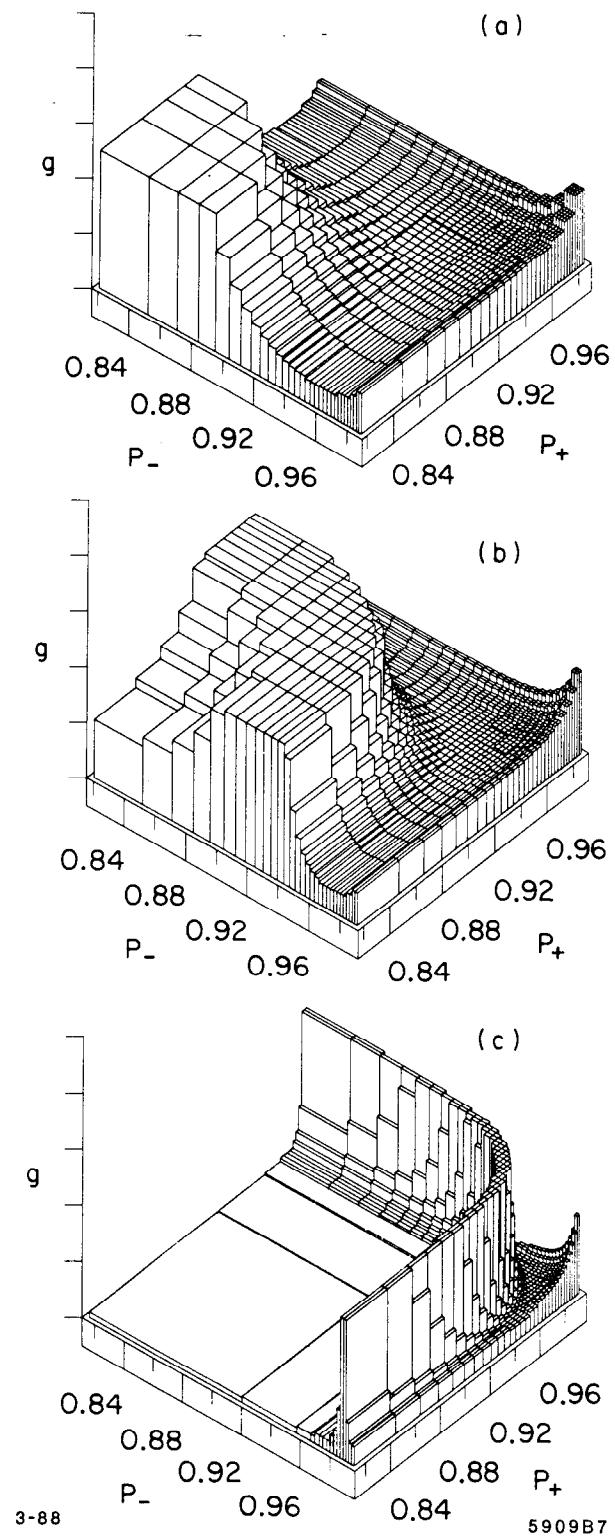
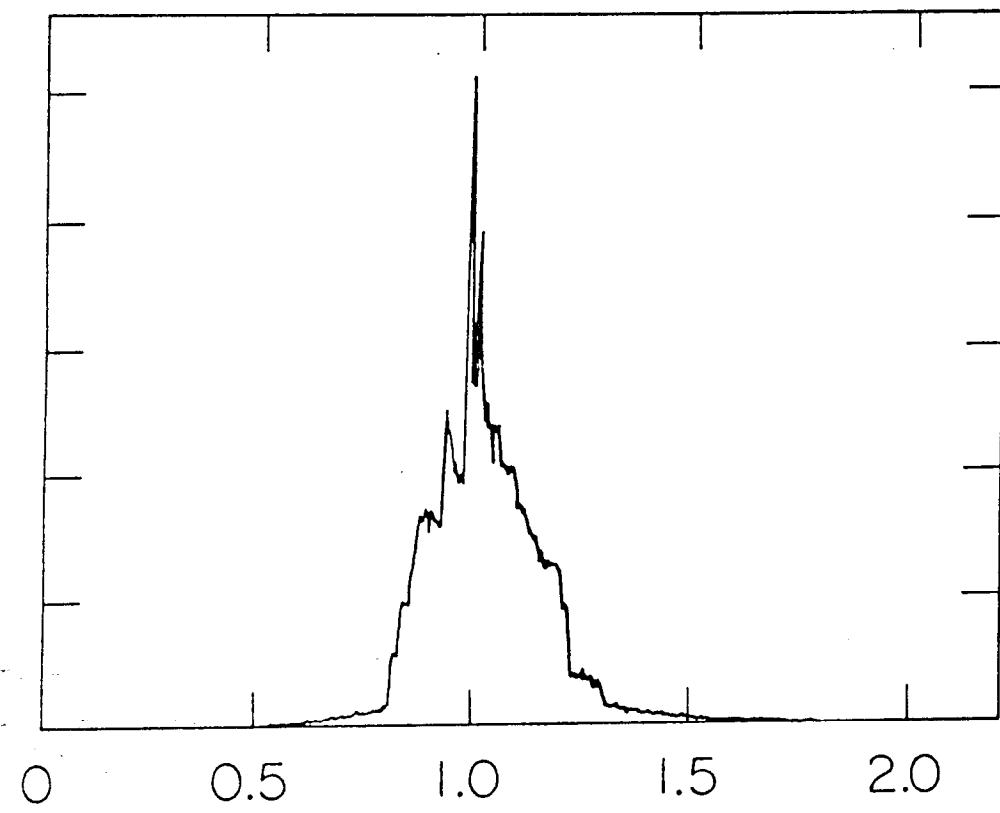


Fig. 4



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Fig. 5

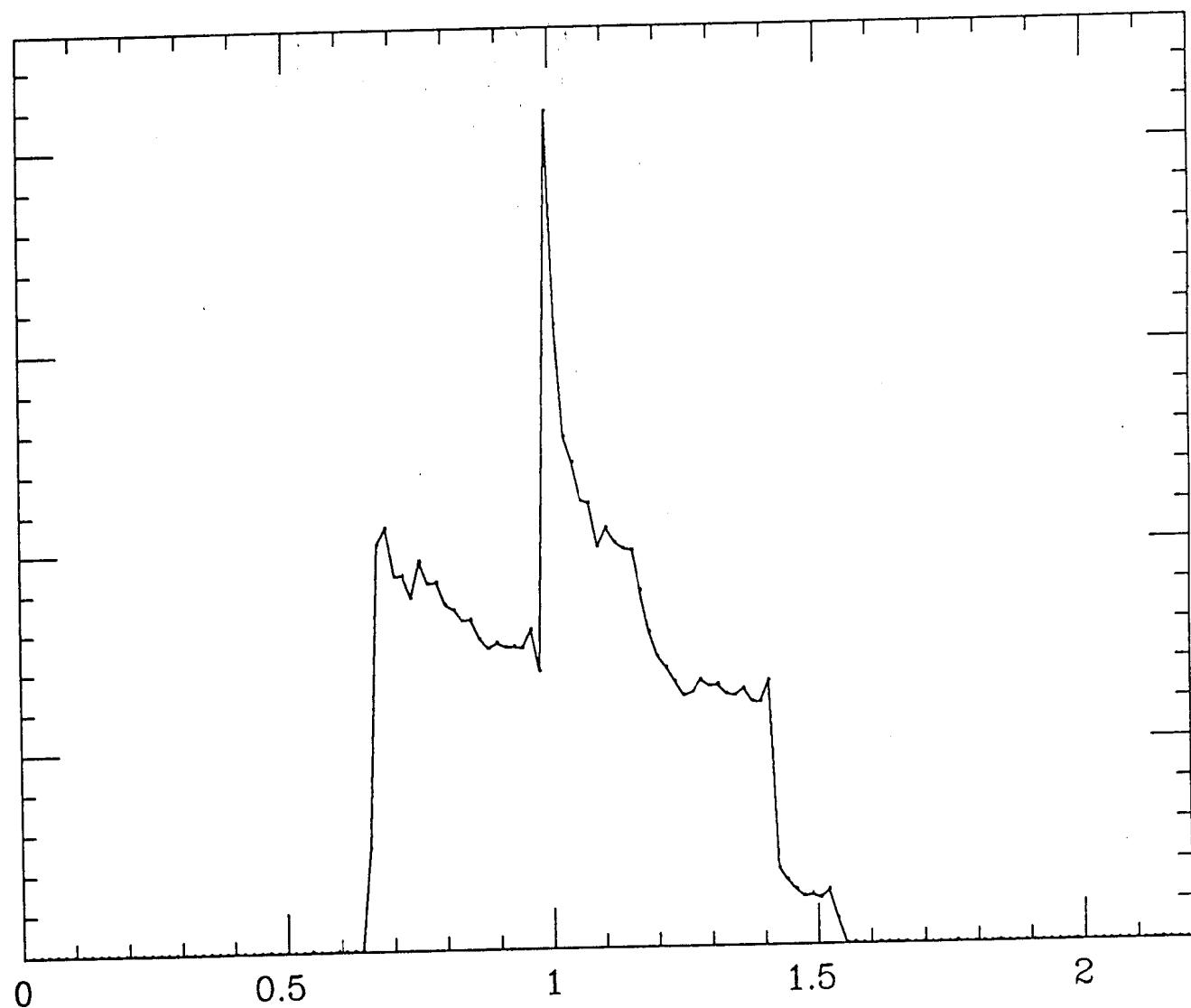


Fig. 6