Structural properties and reaction dynamics of some of the light highly neutron-rich Si, S and Ar isotopes

Mahesh K. Sharma¹, R. N. Panda², Manoj K. Sharma¹, and S. K. Patra³ ¹School of Physics and Materials Science,

Thapar University, Patiala-147 001, India

²Department of Physics, ITER, Siksha O Anusandhan University, Bhubneswar-751 030, India and Institute of Physics, Sachivalaya marg Bhubneshwar-751 005, India.

The availability of radioactive ion beam (RIB) at intermediate energy range open up new opportunities to investigate some exotic phenomenon like neutron and proton halo, Borromean nuclei, neutron skin and bubble effect etc. The measurement of ⁴⁰Mg and ⁴²Al [1] beyond the dripline given by various mass formulae challenge the earlier predictions and still the investigation of dripline is an interesting topic in nuclear physics. Therefore, we investigate the structural properties and reaction dynamics of some of the isotopes of ${}^{32-35}Si$, ${}^{34-37}S$ and ${}^{34-48}Ar$ nuclei, which are well beyond the β - stability line. The bulk properties like binding energy (B.E.), charge radius (r_c) are calculated with well known relativistic mean field (RMF) [2] and newly developed non relativistic Hartree-Fock Formalism (HF) with simple effective interaction (SEI) [3]. The calculated values of B.E and r_c are listed in Table I along with experimental data, both showing nice agreement with each other. For reaction dynamics, we used the well known Glauber approach [4]. The reaction cross sections (σ_R) have been calculated by using expression

$$\sigma_R = 2\pi \int_0^\infty b[1 - T(b)]db, \qquad (1)$$

where 'T(b)' is the Transparency function with impact parameter 'b'. The main ingredient of Glauber model are the densities of projectile and target nuclei. The densities are taken from RMF(NL3) and HF(SEI). The densities of chosen set of nuclei are plotted as a radial distance in Fig.1 using both formalism.

The spherical densities of projectile and target are converted into the Gaussian form in terms of their coefficient c_i 's and a_i 's. These

TABLE I: The calculated values of binding energy (B.E.) in MeV and charge radius (r_c) in fm for projectile and target nuclei from RMF(NL3) and HF(SEI) along with experimental data [5, 6].

| Nuclei | B.E | | | r_c | | |
|--------------------|---------|--------|---------|-------|-------|-------|
| | HF | RMF | Expt. | HF | RMF | Expt |
| | SEI | NL3 | | SEI | NL3 | |
| ^{12}C | 88.422 | 88.23 | 92.160 | 2.436 | 2.364 | 2.47 |
| ³² Si | 267.928 | 267.69 | 271.407 | 3.095 | 3.113 | |
| ³³ Si | 275.953 | 275.97 | 275.915 | 3.095 | 3.133 | |
| ^{34}Si | 283.470 | 283.78 | 283.428 | 3.111 | 3.153 | |
| ^{35}Si | 288.258 | 289.77 | 285.903 | 3.119 | 3.163 | |
| ^{34}S | 286.424 | 286.05 | 291.838 | 3.199 | 3.270 | 3.28 |
| ^{35}S | 296.491 | 296.71 | 298.824 | 3.209 | 3.282 | |
| ^{36}S | 305.978 | 306.52 | 308.714 | 3.219 | 3.293 | 3.29 |
| ^{37}S | 312.676 | 313.58 | 313.017 | 3.224 | 3.299 | |
| $^{34}\mathrm{Ar}$ | 273.357 | 273.76 | 278.719 | 3.295 | 3.387 | 3.365 |
| $^{36}\mathrm{Ar}$ | 300.129 | 300.25 | 306.716 | 3.305 | 3.388 | 3.390 |
| $^{38}\mathrm{Ar}$ | 324.432 | 325.59 | 327.342 | 3.318 | 3.397 | 3.402 |
| $^{40}\mathrm{Ar}$ | 341.555 | 342.31 | 343.810 | 3.326 | 3.401 | 3.427 |
| $^{42}\mathrm{Ar}$ | 357.013 | 357.39 | 359.335 | 3.335 | 3.406 | 3.435 |
| $^{44}\mathrm{Ar}$ | 371.135 | 371.42 | 373.728 | 3.345 | 3.410 | 3.445 |
| $^{46}\mathrm{Ar}$ | 383.672 | 384.57 | 386.927 | 3.373 | 3.410 | 3.437 |
| $^{48}\mathrm{Ar}$ | 391.79 | 394.58 | 396 | 3.393 | 3.437 | |

coefficients are listed in Table II. Other ingredients for reaction cross sections are some energy dependent parameters σ_{NN} , α_{NN} and β_{NN} , which are estimated using Ref. [7]. Figure 2 shows the variation of σ_R as a function of projectile energy (E_{proj}) , for RMF(NL3) and HF(SEI) densities. The σ_R show large magnitude at lower E_{proj} , which start decreasing upto 300 MeV/nucleon. Small enhancement in σ_R is observed up to 300-750 MeV/nucleon and after that, it remains constant. From Fig. 2, we observed that the σ_R values obtained from the HF(SEI) densities are slightly higher

Available online at www.sympnp.org/proceedings



FIG. 1: Radial density plots for $^{32-35}Si$, $^{34-37}S$ and $^{34-48}Ar$ nuclei for (A) HF(SEI) and (B) RMF(NL3).



FIG. 2: Variation of total reaction cross section as function of projectile energy using (A) HF(SEI) and (B) RMF(NL3).

than one obtained using RMF(NL3) densities.

References

- T. Baumann et al., Nature 449, 1022 (2007).
- [2] S. K. Patra and C. R. Praharaj, Phys. Rev. C 44, 2552 (1991).
- [3] B. Behera, X. Viñas, M. Bhuyan, T. R. Routray, B. K. Sharma and S. K. Patra, J. phys. G: Nucl. Part. Phys. (2013, in press).

TABLE II: The Gaussian coefficients c_1 , a_1 , c_2 , a_2 of projectile and target nuclei for RMF(NL3) and HF(SEI) densities.

| Nuclei | | c_1 | a_1 | c_2 | a_2 |
|------------|-----------------------------|-----------|----------|------------|----------|
| ^{12}C | HF(SEI) | -3.79616 | 0.361674 | 3.98071 | 0.345423 |
| | RMF | -0.232654 | 0.638687 | 0.0.517232 | 0.339911 |
| ^{32}Si | HF(SEI) | -2.37277 | 0.202381 | 2.51805 | 0.184376 |
| | RMF | -3.00897 | 0.227291 | 3.13661 | 0.206144 |
| ^{33}Si | HF(SEI) | -2.35386 | 0.196639 | 2.49616 | 0.179151 |
| | RMF | -2.9272 | 0.218848 | 3.05298 | 0.198521 |
| ^{34}Si | HF(SEI) | -2.34203 | 0.191388 | 2.48123 | 0.174374 |
| | RMF | -2.86769 | 0.211811 | 2.99191 | 0.192149 |
| ^{35}Si | HF(SEI) | -2.36933 | 0.187986 | 2.50617 | 0.171288 |
| | RMF | -2.88628 | 0.208281 | 3.00941 | 0.188926 |
| ${}^{34}S$ | HF(SEI) | -1.81296 | 0.181829 | 1.98726 | 0.165612 |
| | RMF | -2.03669 | 0.193789 | 2.20712 | 0.175664 |
| ${}^{35}S$ | HF(SEI) | -1.78947 | 0.176802 | 1.96139 | 0.161028 |
| | RMF | -2.11099 | 0.189798 | 2.27085 | 0.172077 |
| ${}^{36}S$ | HF(SEI) | -1.76837 | 0.172069 | 1.93796 | 0.156714 |
| | RMF | -2.13714 | 0.185085 | 2.28947 | 0.167819 |
| ${}^{37}S$ | HF(SEI) | -1.79518 | 0.169348 | 1.96303 | 0.154256 |
| | RMF | -2.21565 | 0.183299 | 2.3634 | 0.166219 |
| ^{34}Ar | HF(SEI) | -1.93044 | 0.183573 | 2.09551 | 0.167211 |
| | RMF | -1.84812 | 0.18999 | 2.02988 | 0.172144 |
| ^{36}Ar | HF(SEI) | -1.83132 | 0.172884 | 1.99595 | 0.157479 |
| | RMF | -1.79242 | 0.177609 | 1.96319 | 0.160928 |
| ^{38}Ar | HF(SEI) | -1.79848 | 0.164912 | 1.96024 | 0.150194 |
| | RMF | -1.82537 | 0.169399 | 1.98489 | 0.153507 |
| ^{40}Ar | HF(SEI) | -1.85865 | 0.160185 | 2.01657 | 0.145936 |
| | RMF | -1.99749 | 0.166945 | 2.14665 | 0.151332 |
| ^{42}Ar | HF(SEI) | -1.91109 | 0.155909 | 2.06565 | 0.142049 |
| | RMF | -2.26358 | 0.166337 | 2.39771 | 0.150869 |
| ^{44}Ar | HF(SEI) | -1.96876 | 0.151999 | 2.11994 | 0.138517 |
| | RMF | -2.68801 | 0.168372 | 2.79633 | 0.152825 |
| ^{46}Ar | $\mathrm{HF}(\mathrm{SEI})$ | -2.71167 | 0.158333 | 2.8162 | 0.144599 |
| | RMF | -3.13997 | 0.170354 | 3.22018 | 0.154714 |
| ^{48}Ar | $\mathrm{HF}(\mathrm{SEI})$ | -2.68582 | 0.155294 | 2.80113 | 0.141718 |
| | RMF | -3.15705 | 0.167492 | 3.24426 | 0.152135 |

- [4] R. J. Glauber, Phys. Rev. **100**, 1 (1955).
- [5] G. Audi, A. H. Wapstra and C. Thibault, Nucl. Phys. A 729, 337 (2003).
- [6] I. Angeli, K. P. Marinova, Atomic Data and Nuclear Data Table 99, 69-95 (2013).
- [7] W. Horiuchi, Y. Suzuki, B. Abu Ibrahim and A. Kohama, Phys. Rev. C 75, 044607 (2007).