

RECENT CLEO RESULTS ON CHARM PHYSICS

E. I. SHIBATA (CLEO Collaboration)
 Department of Physics, Purdue University
 West Lafayette, IN 47907, U. S. A.

ABSTRACT

A summary of recent charm physics results from CLEO 1.5 is presented. Branching fractions for D^0 decays to $K+K^-$, K^0K^0 , $\mathbf{TT}+\mathbf{TT}'$, \mathbf{TTV} , and some other modes involving a w^0 or an $r^?$ are given. Using the observation of $Df \rightarrow \langle j \rangle l + \nu$, the absolute branching fraction for $Df \rightarrow \langle j \rangle n +$ has been derived. A_c results include branching fractions into pK^-x , pK^0 , $pR^{\wedge\wedge}$, $A^?r+$, and $Air^?w+$, and the A_c decay asymmetry parameter.

Introduction

The charm physics results reported in this paper are from e^+e^- annihilation data collected by the CLEO 1.5 detector¹ at the Cornell Electron Storage Ring (CESR) in Ithaca, New York. The data set corresponds to integrated luminosities of 101 pb^{-1} below the T(4S) and 212 pb^{-1} at the T(4S) taken in 1987 and 116 pb^{-1} taken at the T(5S) in early 1988. About $\sim 1,100,000$ hadronic events from the continuum are contained in this data set. Throughout this paper charge conjugate states are implied.

D^0 decays

While many of the theoretical calculations for rates of two-body, non-leptonic decays of the D^0 are in agreement with experimental measurements, $D^0 \rightarrow KK$ and $D^0 \rightarrow ww$ present problems. The current world average² for the ratio of branching fractions $B(D^0 \rightarrow K+K^-)/B(D^0 \rightarrow *+*^-)$ is 3.9 ± 1.2 . This is not easily reconciled with theoretical expectations³ which range from 1 to 1.4. In lowest order the process $D^0 \rightarrow K^0K^0$ proceeds through two VP-exchange diagrams whose sum cancels in the limit of exact $SU(2)$ flavor symmetry, so $B(D^0 \rightarrow K^0K^0)$ is predicted to be small⁴ ($\sim 10^{-4}$ or less) in a simple quark picture.

Using $19^{*+} \rightarrow D^0\mathbf{TT}+$ events, selected by requiring $|AAf - 145.45 \text{ MeV}/c^2| < 2.4 \text{ MeV}/c^2$, where $AM = M(D^{*+}) - M(D^0)$ and $x(D^{*+}) = p/p_{max} > 0.5$, the $K+K^-$ and $\pi^+\pi^-$ invariant mass distributions shown in Figs. 1(a) and 1(b)

were obtained. The peaks centered at an invariant mass of $1.865 \text{ GeV}/c^2$ are from $D^0 \rightarrow \mathbf{JfT}+\mathbf{iT}$ (249 ± 21 events) and $D^0 \rightarrow \mathbf{TT}+\mathbf{TT}$ (110 ± 15 events), respectively. The other structures are due to reflections from the copious $D^0 \rightarrow K^{\wedge}w^+$ and $D^0 \rightarrow$

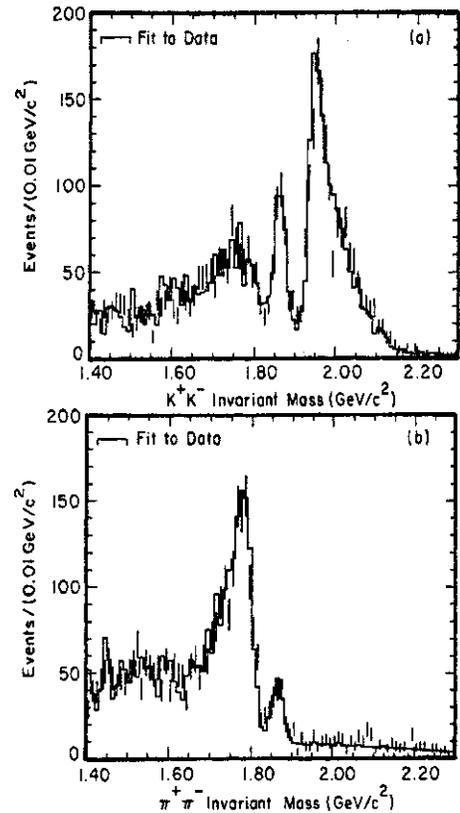


Fig. 1. Invariant mass distributions for (a) $K+K^-$ and (b) $\pi^+\pi^-$. The fit to the data is by the sum of a Monte Carlo simulated background from D^0 decays, a polynomial background, and a Gaussian signal.

K^0 decay modes. Normalizing to the decay channel $D^0 \rightarrow K^+ K^-$ and correcting for efficiencies, we find $B(D^0 \rightarrow K^+ K^-) = (0.49 \pm 0.04 \pm 0.03 \pm 0.06)\%$ and $B(D^0 \rightarrow \tau^+ \tau^-) = (0.21 \pm 0.03 \pm 0.02 \pm 0.03)\%$, where the third error is due to the uncertainty in $B(D^0 \rightarrow K^+ K^-) = (4.2 \pm 0.6)\%$. Thus, the ratio of branching fractions $B(D^0 \rightarrow K^+ K^-)/B(D^0 \rightarrow \tau^+ \tau^-)$ is $2.35 \pm 0.37 \pm 0.28$, lower than the current world average but higher than theoretical expectations.

$D^0 \rightarrow K^0 K^0$ candidates were selected from $D^{*+} \rightarrow D^0 \pi^+$ events by requiring that $|M(D^0 \rightarrow K^0 K^0) - 145.45 \text{ MeV}/c^2| < 1.2 \text{ MeV}/c^2$, $M(D^0 \rightarrow K^0 K^0)$ to be within $12.5 \text{ MeV}/c^2$ of $M(K^0)$ and $x(D^{*+}) > 0.5$. We observe 5 events with masses consistent with D^0 decay. From Monte Carlo simulations the background is estimated to be 0.3 events. In order to reduce the systematic error in the determination of the $D^0 \rightarrow K^0 K^0$ branching fraction, we normalized to the decay channel $D^0 \rightarrow K^+ K^-$. Using the branching ratio¹⁵¹ $D^0 \rightarrow K^+ K^- = (6.4 \pm 1.1)\%$, we find $B(D^0 \rightarrow K^0 K^0) = (0.13^{+0.07}_{-0.02})\%$, where the systematic error is dominated by the uncertainty in $B(D^0 \rightarrow K^+ K^-)$. This result is consistent with Pham's calculation¹ based on non-perturbative hadronic final state interactions, in which he obtained $B(D^0 \rightarrow K^0 K^0) \ll B(D^0 \rightarrow K^+ K^-) \ll 0.25\%$. Here we have used our value for $B(D^0 \rightarrow K^+ K^-)$ given above.

D^0 decays involving a τ^0 or an n

Branching fractions for D^0 decays involving a τ^0 or an n are summarized in the Table 1 below. Also shown are theoretical predictions by Bauer, Steck, and Wirbel¹⁵¹ (BSW) and Blok and Shifman¹⁵¹ (BS). Our measurements are in good agreement with the BSW predictions, but are somewhat higher than the BS predictions.

Table 1. D^0 branching fractions for decays involving a τ^0 or an n .

Mode	CLEO	BSW	BS
$K^+ \pi^+ \pi^0$	$11.5 \pm 0.6 \pm 2.1$		
$\bar{K}^0 \pi^0$	$2.3 \pm 0.4 \pm 0.5$	2.5	1.5
$K^+ \pi^+ \pi^- \pi^+ \pi^0$	$5.0 \pm 0.7^{+1.3}_{-1.0}$		
$\bar{K}^0 \omega$	$3.4 \pm 0.9 \pm 1.0$	2.7	1.5
$\bar{K}^0 \eta$	$2.3^{+0.7}_{-1.1}$	2.5	0.3
$\pi^0 \pi^0$	< 0.46 (90% C.L.)		

$D^0 \rightarrow \tau^+ \tau^-$ and $D^0 \rightarrow n \tau^+$.

Through the observation of $D^0 \rightarrow \tau^+ \nu$, we have made a determination of the absolute branching fraction for $D^0 \rightarrow \tau^+ \nu$. It is found that the cuts $p(\tau^+) > 2 \text{ GeV}/c$ and $p(\nu) > 1 \text{ GeV}/c$ isolate $D^0 \rightarrow \tau^+ \nu$ events. After lepton fake and BE background subtractions there are 37.4 ± 9.0 and 17.0 ± 6.4 events. There are 400 ± 27 $D^0 \rightarrow \tau^+ \nu$ events. Averaging the $D^0 \rightarrow \tau^+ \nu$ and $D^0 \rightarrow \tau^+ \nu$ data samples and correcting for efficiencies, we find

$$B(D^0 \rightarrow \tau^+ \nu)/B(D^0 \rightarrow K^+ K^-) = 0.49 \pm 0.10^{(stat)} \pm 0.05^{(sys)}.$$

The value for the $D^0 \rightarrow \tau^+ \nu$ branching fraction is derived from the following relation:

$$B(D^0 \rightarrow \tau^+ \nu) = \frac{f_{\tau^+ \nu}}{f_{K^+ K^-}} \cdot B(D^0 \rightarrow K^+ K^-) \cdot \frac{\tau_{\tau^+ \nu}}{\tau_{K^+ K^-}}.$$

The factor 0.8 is the average of two predictions,¹⁵¹ and the error reflects a large range of possible differences in form factor. The measured branching fraction¹⁵¹ for $B(D^0 \rightarrow K^+ K^-)$ is $(4.5 \pm 0.7 \pm 0.5)\%$ and the ratio of D_s and D^* lifetimes is 0.42 ± 0.03 . The resulting estimate for $B(D^0 \rightarrow \tau^+ \nu)$ is $(1.50 \pm 0.31)\%$. Thus, $B(D^0 \rightarrow \tau^+ \nu) = (3.1 \pm 0.61^{(stat)} \pm 0.6)^{\%}$, where the first error is statistical, the second is systematic, and the third is also systematic and arises from the uncertainty on the predicted value of $B(D^0 \rightarrow \tau^+ \nu)$. This value for $B(D^0 \rightarrow \tau^+ \nu)$ can be compared with the Mark III upper limit¹⁵¹ of 4.1% and the E691 lower limit¹⁵¹ of 3.4%

A. branching fractions and decay asymmetry

A.1 branching fractions

Absolute branching fractions for several A₁ decay modes are shown in Table 2 along with the values given by the Particle Data Group³ (PDG). The CLEO numbers are based on $B(K_c^- \rightarrow p \bar{K}^- n^+) = (4.3 \pm 1.4)\%$, which is a weighted average¹⁵¹ of CLEO and ARGUS estimates.

Table 2. A₁ branching fractions.

Decay mode	CLEO	PDG
$p \bar{K}^- \pi^+$	4.3 ± 1.4	2.8 ± 0.8
$p \bar{K}^0$	2.1 ± 0.7	1.6 ± 0.6
$p \bar{K}^0 \pi^- \pi^+$	1.8 ± 0.8	8.1 ± 3.5
$\Lambda \pi^+$	0.7 ± 0.3	seen
$\Lambda \pi^+ \pi^- \pi^+$	2.8 ± 1.0	1.9 ± 0.7

A_1 decay asymmetry parameter and A_1 polarization

Violation of parity conservation in the weak decays of charmed baryons is expected. The decay $A^+ \rightarrow A^0 \pi^+$ is analogous to the decay $A \rightarrow p \pi^-$, for which the parity-violating asymmetry decay parameter has been measured to be $\langle A \rangle = 0.642 \pm 0.013$. The form of the angular distribution of the proton in the decay $A^+ \rightarrow A^0 \pi^+$, where $A \rightarrow p \pi^-$, is given by $dN/d\cos\theta_1 = |(1 + \langle A \rangle \cos\theta_1)|^2$, where θ_1 is the angle between the A_1 direction in the A_1 rest frame and the decay proton's line of flight in the A_1 rest frame. The fit to the CLEO data is shown in Fig. 2 and gives $\langle A \rangle = -1.01 \pm 0.14$, constraining $\langle A \rangle$ to physical values, indicating that parity conservation is violated in the weak decay $A^+ \rightarrow A^0 \pi^+$ as is expected.

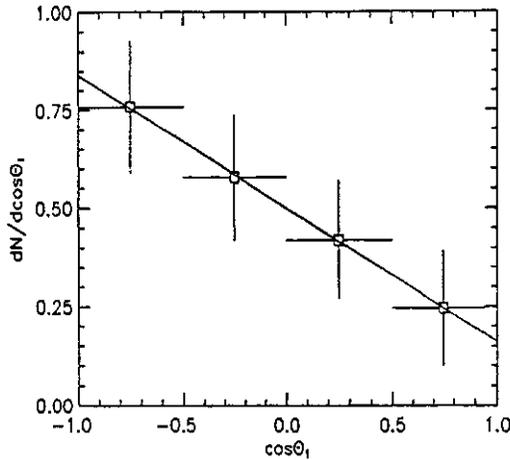


Fig. 2. Angular distribution of the decay proton in the A_1 rest frame. The slope of the distribution is $\langle A \rangle$. The fit line has a slope of -0.34 ± 0.14 .

Parity conservation in electromagnetic annihilation requires A_1 polarization, if it exists, to be normal to the production plane. In addition, the polarization must be the same for particle and antiparticle states since C is a conserved quantum number for A_1 production. We define the normal to the production plane as $\mathbf{n} = \mathbf{p}^+ \times \mathbf{e}^+$, the cross product of the A_1 momentum vector and the direction of the positron beam. In the A^+ rest frame the angular distribution of the A relative to \mathbf{n} has the form $dN/d\cos\theta_2 = |(1 + \langle A \rangle \cos\theta_2)|^2$, where P is the polarization and θ_2 is the angle between \mathbf{n} and the A direction in the A_1 rest frame. Since $\langle A \rangle = -\langle A \rangle$, subtracting the l^+ distribution from the l^- distribution yields $dN/d\cos\theta_2 = +\langle A \rangle \cos\theta_2$. The fit to this distribution, shown in Fig. 3, gives $P = -0.2 \pm 0.2$, assuming that $\langle A \rangle = -1.0$. Thus, we see no evidence for the production of polarized A_1 .

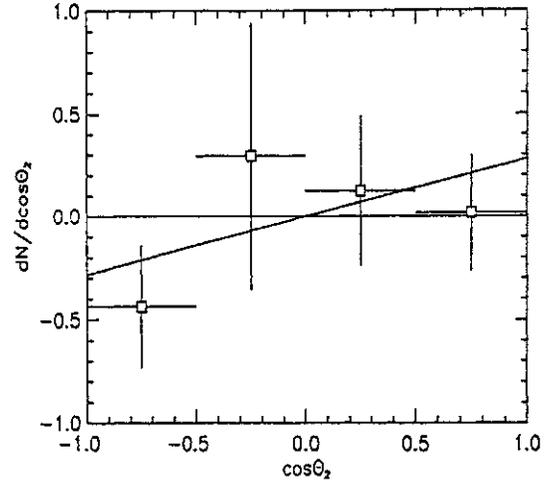


Fig. 3. The angular distribution of the A_1 relative to \mathbf{n} . The slope of the distribution is $+P\langle A \rangle$. The fit line has a slope of $+0.24 \pm 0.24$.

Acknowledgements

The author wishes to thank his CLEO collaborators for their many contributions and helpful suggestions and the CESR staff for their efforts. This work was supported in part by the U. S. Department of Energy under contract DE-AC02-76ER01428.

References

1. D. Andrews et al, *Nuci Inst. and Meth.* 211, 47 (1983); S. Behrends et al., *Phys. Rev.* D31, 2161 (1985); D. Cassel et al, *Nuci Inst and Meth.* A252, 325 (1986).
2. Particle Data Group, *Phys. Lett* 239, 1 (1990).
3. I.I. Bigi, in *Heavy Quark Physics*, edited by P. S. Drell and D. L. Rubin, AIP Conference Proceedings No. 196 (American Inst. of Phys., New York, 1989), p. 18, and references therein.
4. X.-Y. Pham, *Phys. Lett* 193B, 331 (1987).
5. D. Hitlin, *Nuci. Phys. B (Proc. Suppl.)* 3, 179 (1988).
6. M. Bauer, B. Stech, and M. Wirbel, *Z. Phys.* C54, 103 (1987).
7. B. Yu. Blok and M. A. Shifman, *Sov. J. Nuci. Phys.* 45, 522 (1987).
8. J. Alexander et al, *Phys. Rev. Lett* 85, 1531 (1990).
9. N. Isgur et al. (private communication) and N. Isgur et al. *Phys. Rev.* D39, 799 (1989) predict 0.78 for the ratio $T(D^+ \rightarrow \langle l^+ \nu \rangle) / T(D^+ \rightarrow R^{*+} l^+ \nu)$. M. Wirbel (private communication) predicts 0.83; see M. Wirbel et al, *Z. Phys* C29, 269 (1985).
10. J. C. Anjos et al, *Phys. Rev. Lett* 62, 722 (1989).
11. J. Adler et al, *Phys. Rev. Lett* 64, 169 (1990).
12. J. C. Anjos et al, *Phys. Rev. Lett* 64, 2885 (1990).
13. S. Stone, in *Session Summary: Heavy Quark Decay, XII International Workshop on Weak Interactions and Neutrinos, 1989, Sea of Galilee, Israel.*

DISCUSSION

- Q.* A. N. Kamalff/niv. *Alberta*) : I am a little surprised that you said that the theoretical expectation for the ratio $B(D^0 \rightarrow K+K^-)/B(D^0 \rightarrow \pi^+\pi^-)$ is 1 to 1.4. In fact, it is easy to get a value of 2, and if one is prepared to play with QCD coefficients a_1 and a_2 of Bauer, Stech and Wirbel, one can get up to 3 for this ratio, putting in final state interactions.
- A.* E. Shibata: That is good news. It shows that final state interactions are important.